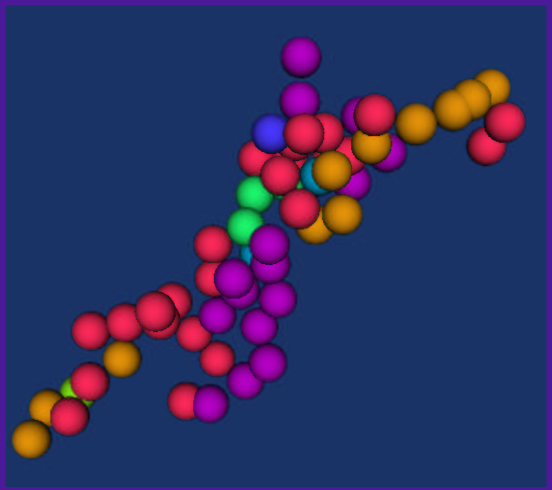
Self-Organized Criticality In a Glass

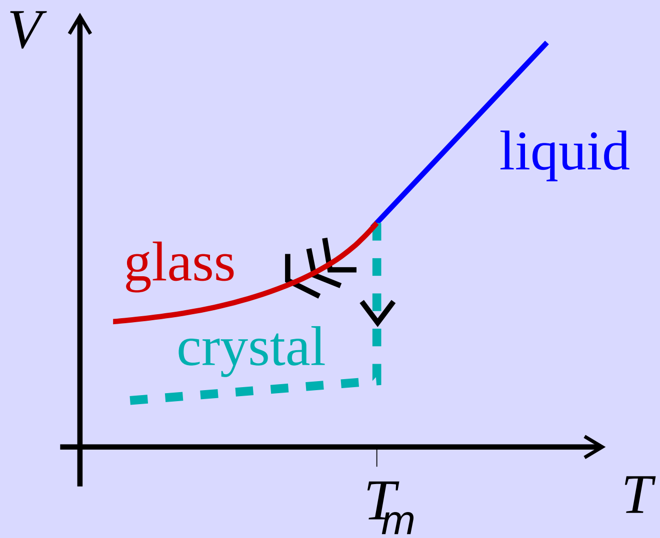
Katharina Vollmayr-Lee, Bucknell University

January 3, 2008



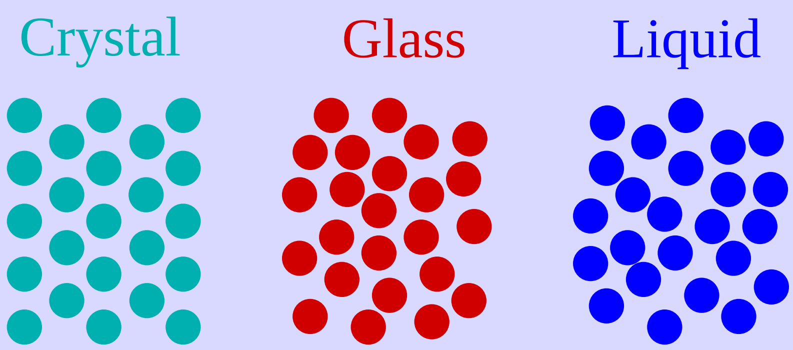
Thanks: E. A. Baker, A. Zippelius, K. Binder, and J. Horbach

Introduction



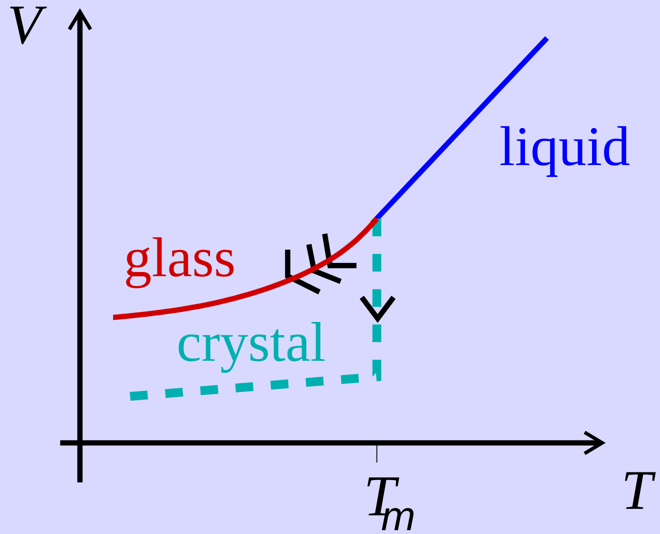
Glass:

→ system falls out of equilibrium



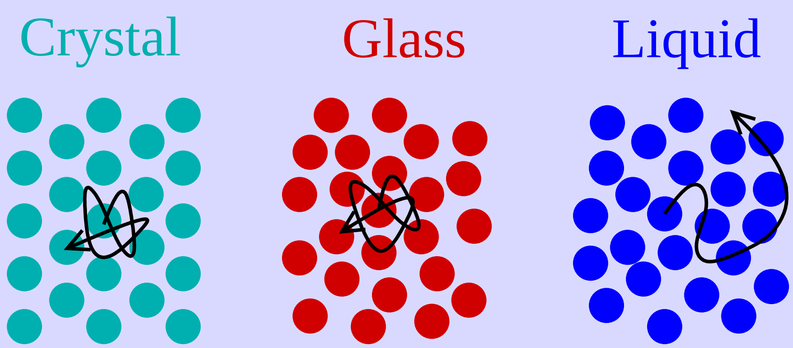
Structure: discordered

Introduction



Glass:

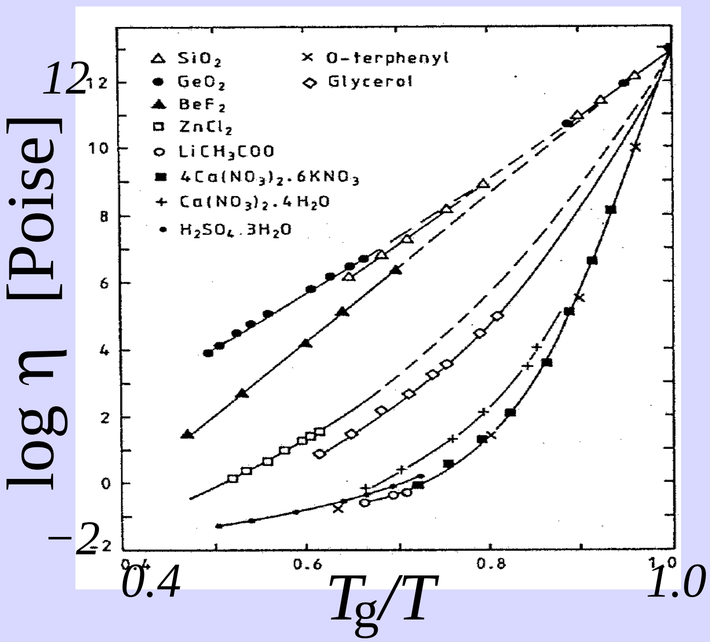
⟶ system falls out of equilibrium



Structure: discordered

Dynamics: frozen in

Introduction



[C.A. Angell and W. Sichina, Ann. NY Acad.Sci. 279, 53 (1976)]

Dynamics:

⟶ slowing down of many decades

Model

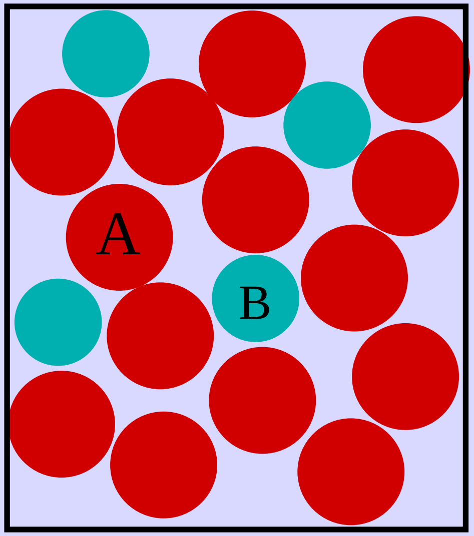
Binary Lennard-Jones System

$$V_{\alpha\beta}(r) = 4\,\epsilon_{\alpha\beta}\left(\left(\frac{\sigma_{\alpha\beta}}{r}\right)^{12} - \left(\frac{\sigma_{\alpha\beta}}{r}\right)^{6}\right)$$

$\sigma_{\mathsf{AA}} = 1.0 \quad \sigma_{\mathsf{AB}} = 0.8 \quad \sigma_{\mathsf{BB}} = 0.88$

$\epsilon_{\mathsf{AA}} = 1.0 \quad \epsilon_{\mathsf{AB}} = 1.5 \quad \epsilon_{\mathsf{BB}} = 0.5$

[W. Kob and H.C. Andersen, PRL 73, 1376 (1994)]



800 A and 200 B

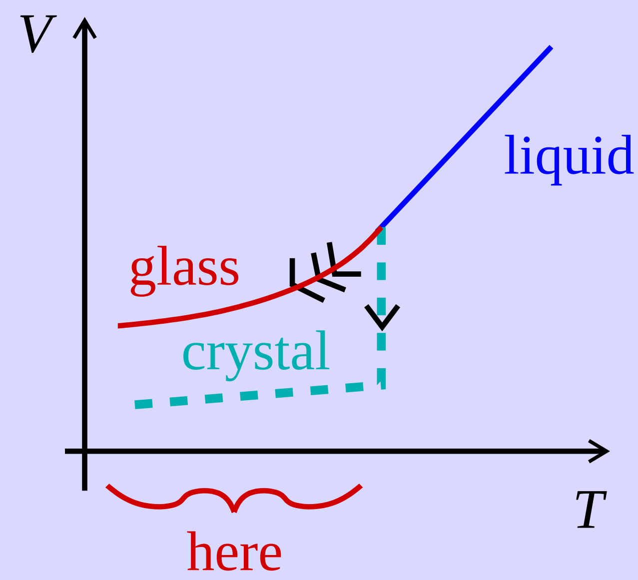
Simulations

Molecular Dynamics Simulations

Velocity Verlet

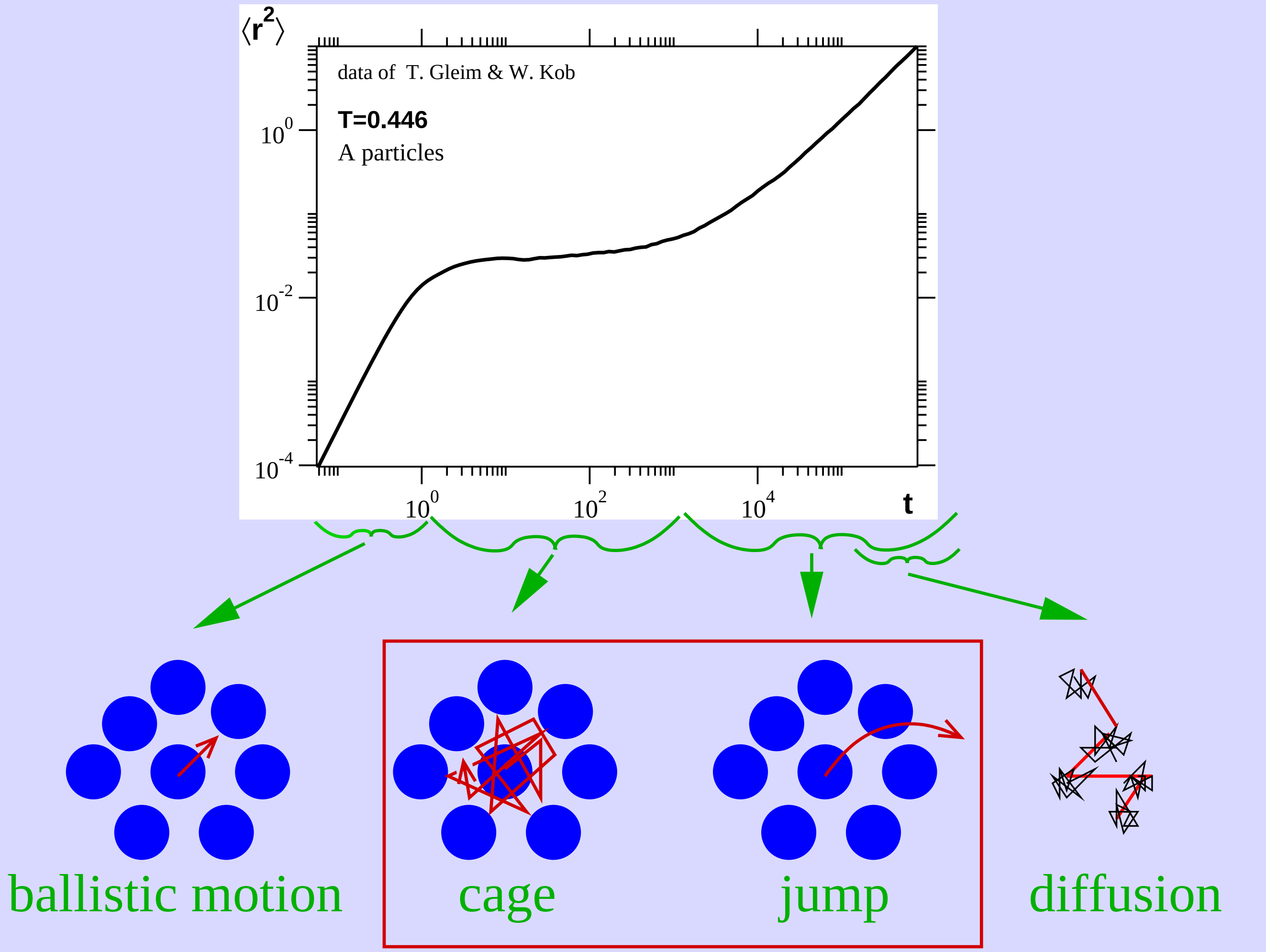
below glass transition:

$T = 0.15 - 0.43 \qquad T_{\mathsf{c}} = 0.435$



Cage-Picture

Mean-Squared Displacement: $\langle r^2 \rangle(t) = \left\langle \frac{1}{N} \sum_{i=1}^{N} \left(\underline{r}_i(t) - \underline{r}_i(0) \right)^2 \right\rangle$

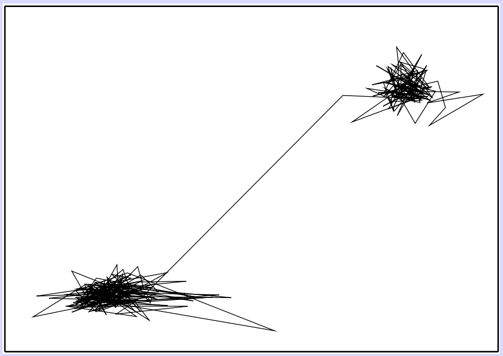


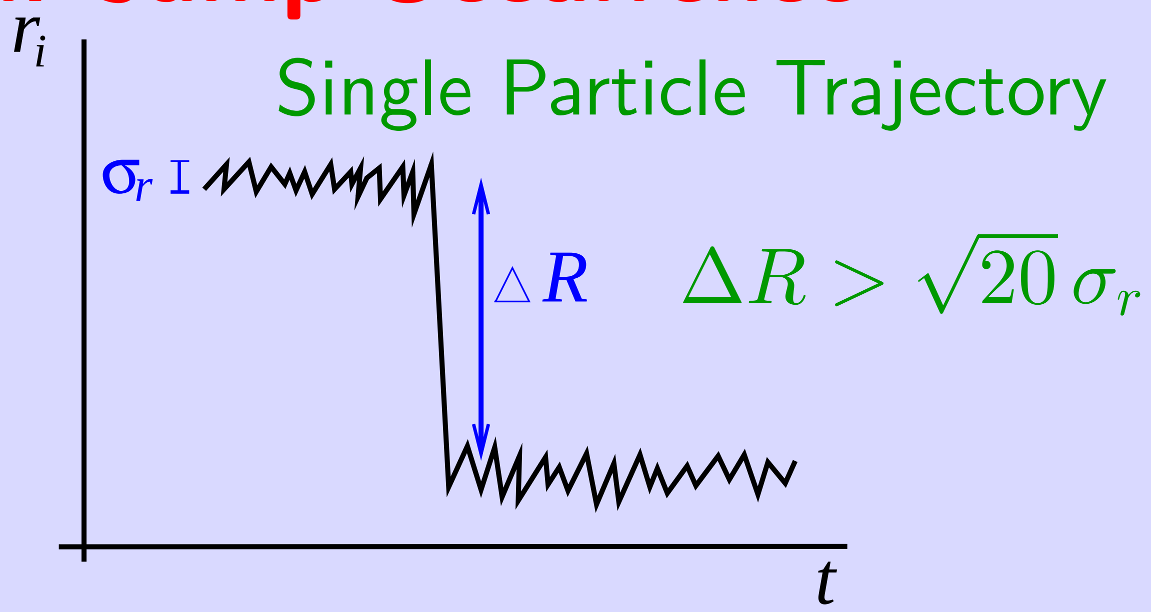
$\langle r^2 \rangle \sim t^2$

here

$\langle r^2 \rangle \sim t$

Definition: Jump Occurrence





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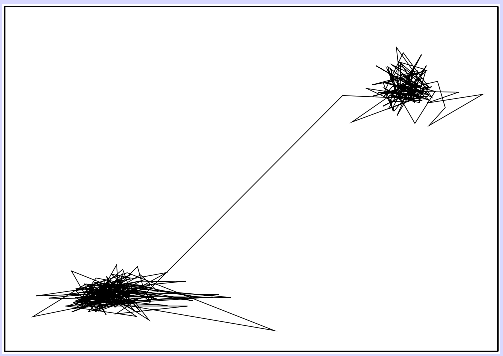
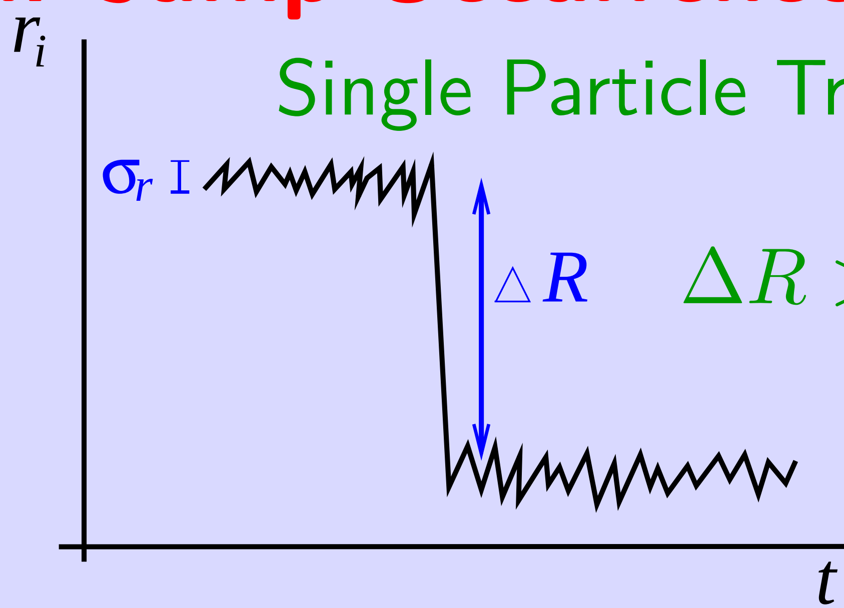
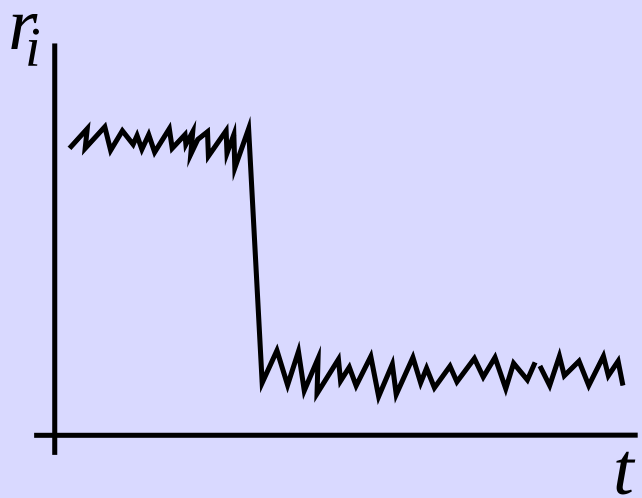




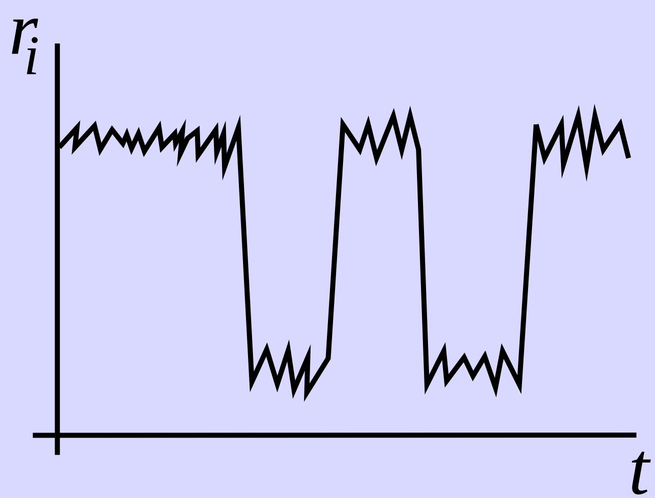
$$\Delta R > \sqrt{20}\,\sigma_r$$

Definition: Jump Type

Irreversible Jump



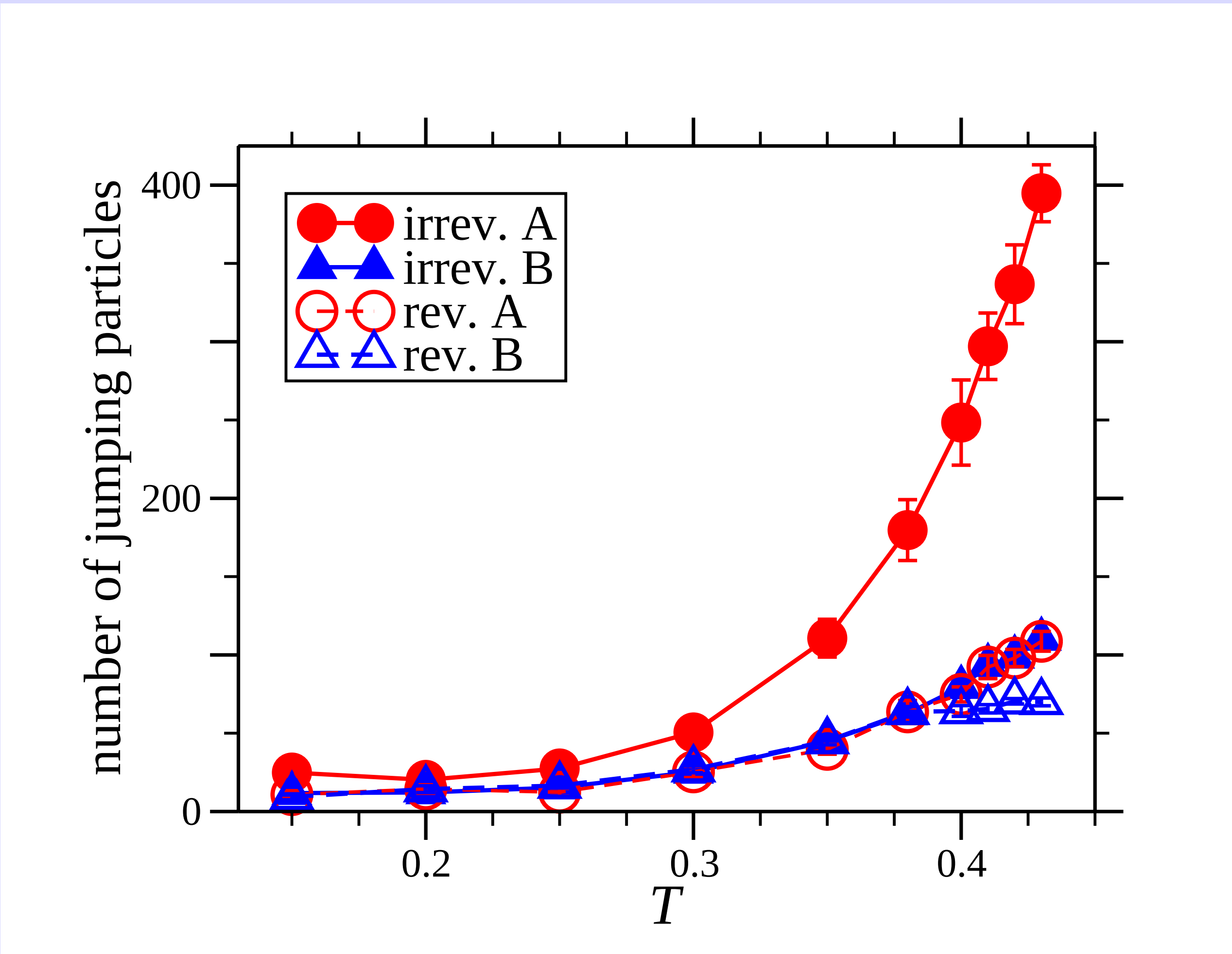
Reversible Jump



Outline

- Jump Statistics
- Correlated Single Particle Jumps
- History Dependence
- Summary

Number of Jumping Particles



$\Longrightarrow$ increasing with increasing T

$\Longrightarrow$ both A & B particles jump

$\Longrightarrow$ irrev. & reversible jumps at all temperatures T

Fraction of Irreversibly Jumping Particles

$$\text{fraction of irrev. jumpers} = \frac{\text{number of irrev. jump. part.}}{\text{number of jump. part.}}$$

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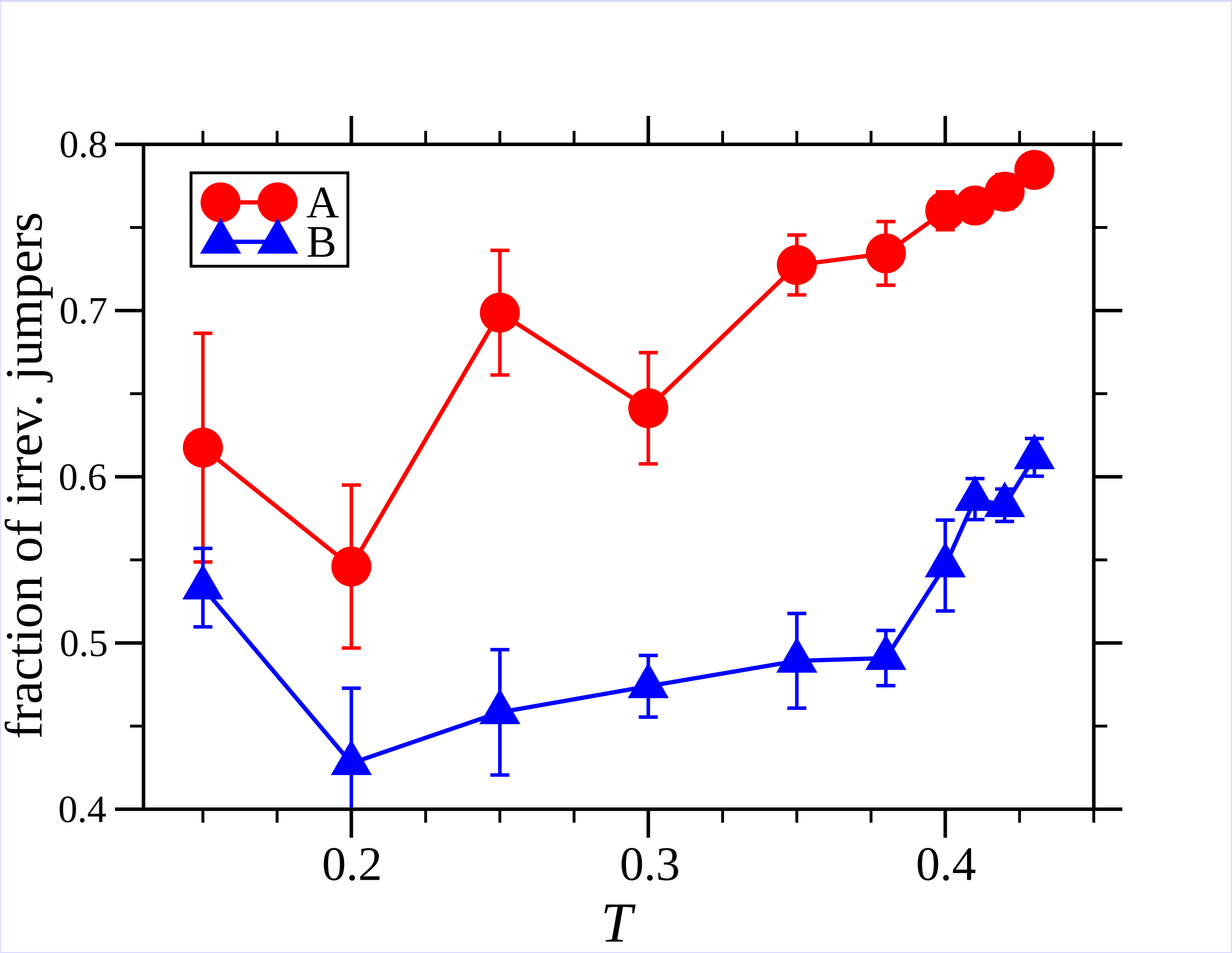
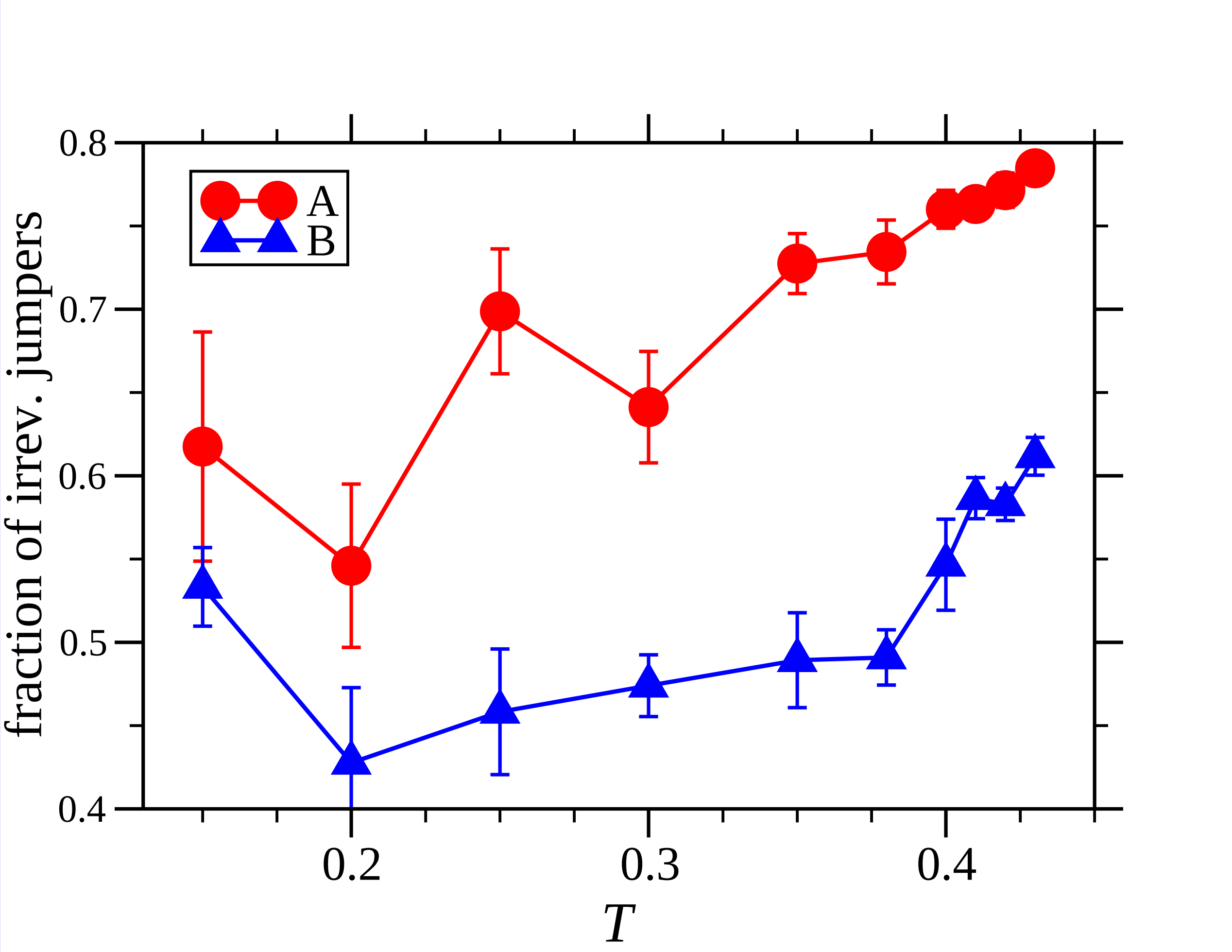


$\Longrightarrow$ fraction of irrev. jumpers increases with increasing T

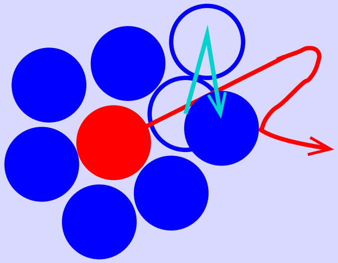
Fraction of Irreversibly Jumping Particles

$$\text{fraction of irrev. jumpers} = \frac{\text{number of irrev. jump. part.}}{\text{number of jump. part.}}$$

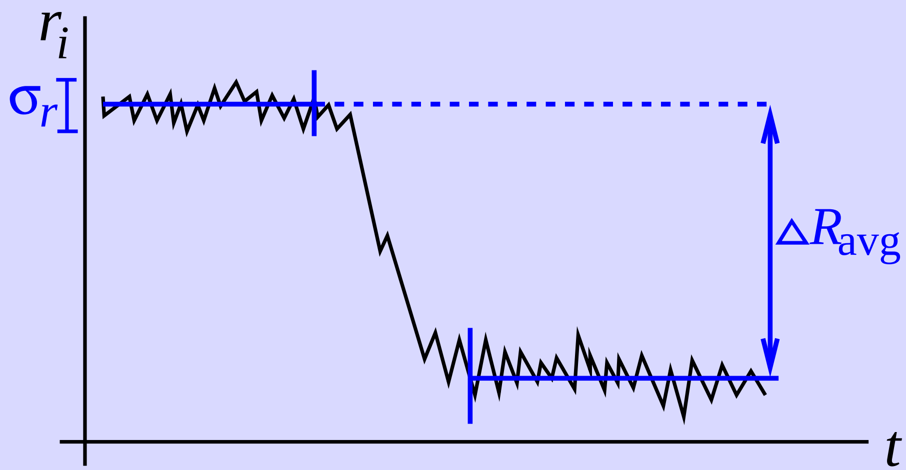


$\Longrightarrow$ fraction of irrev. jumpers increases with increasing T

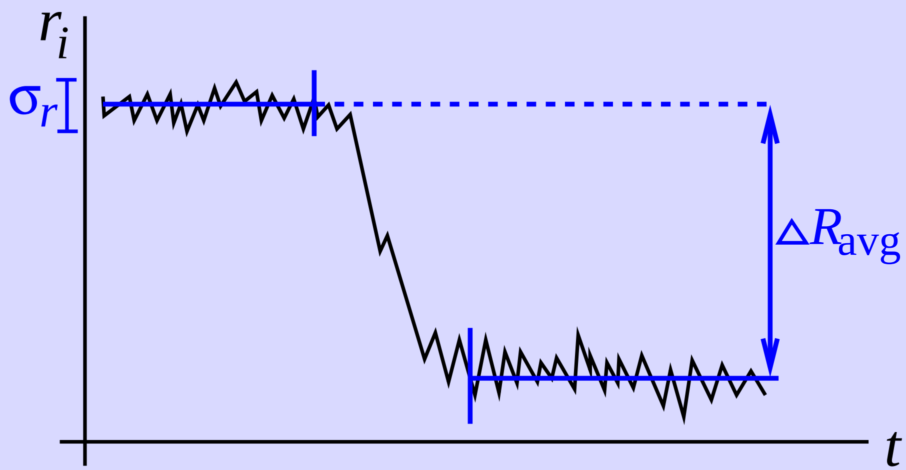
interpretation:
door closing

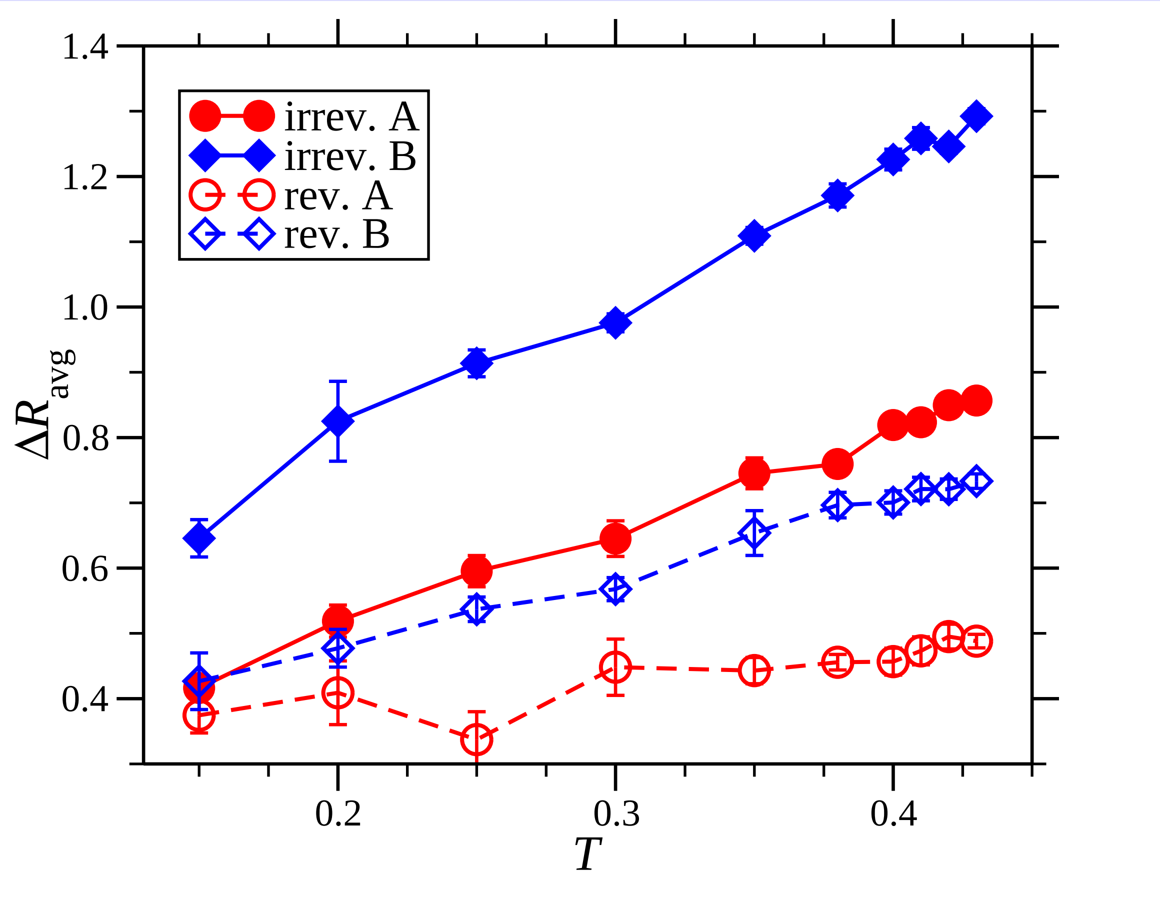


Jump Size



Jump Size



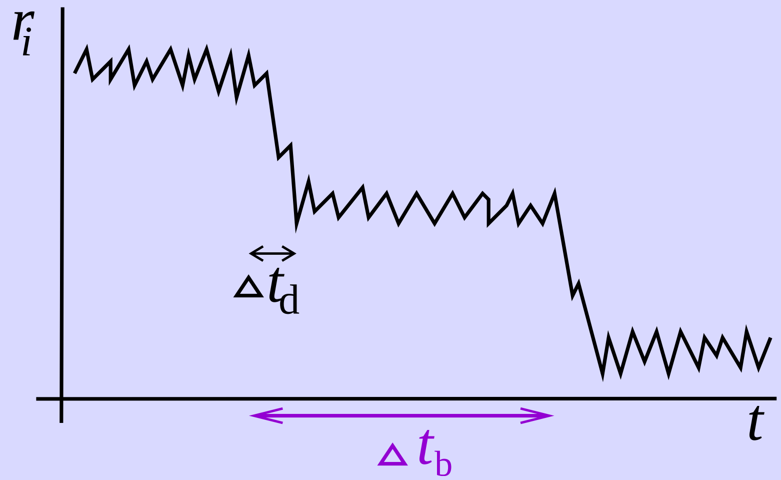


$\Longrightarrow$ increasing with increasing T

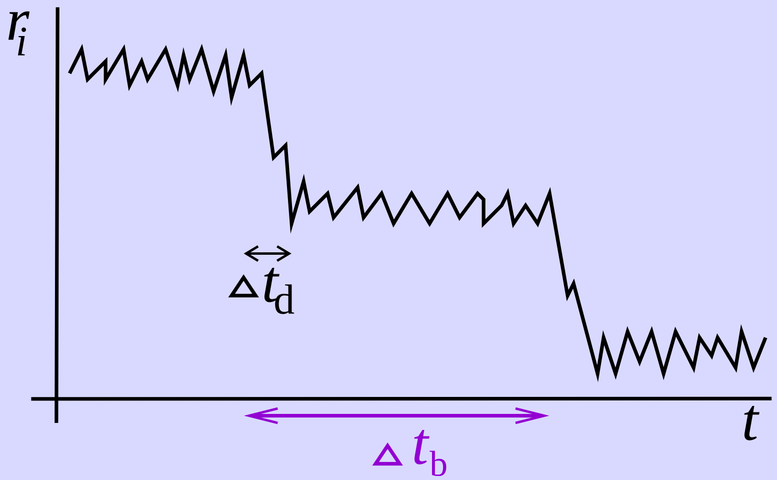
$\Longrightarrow$ (smaller) B-particles jump farther

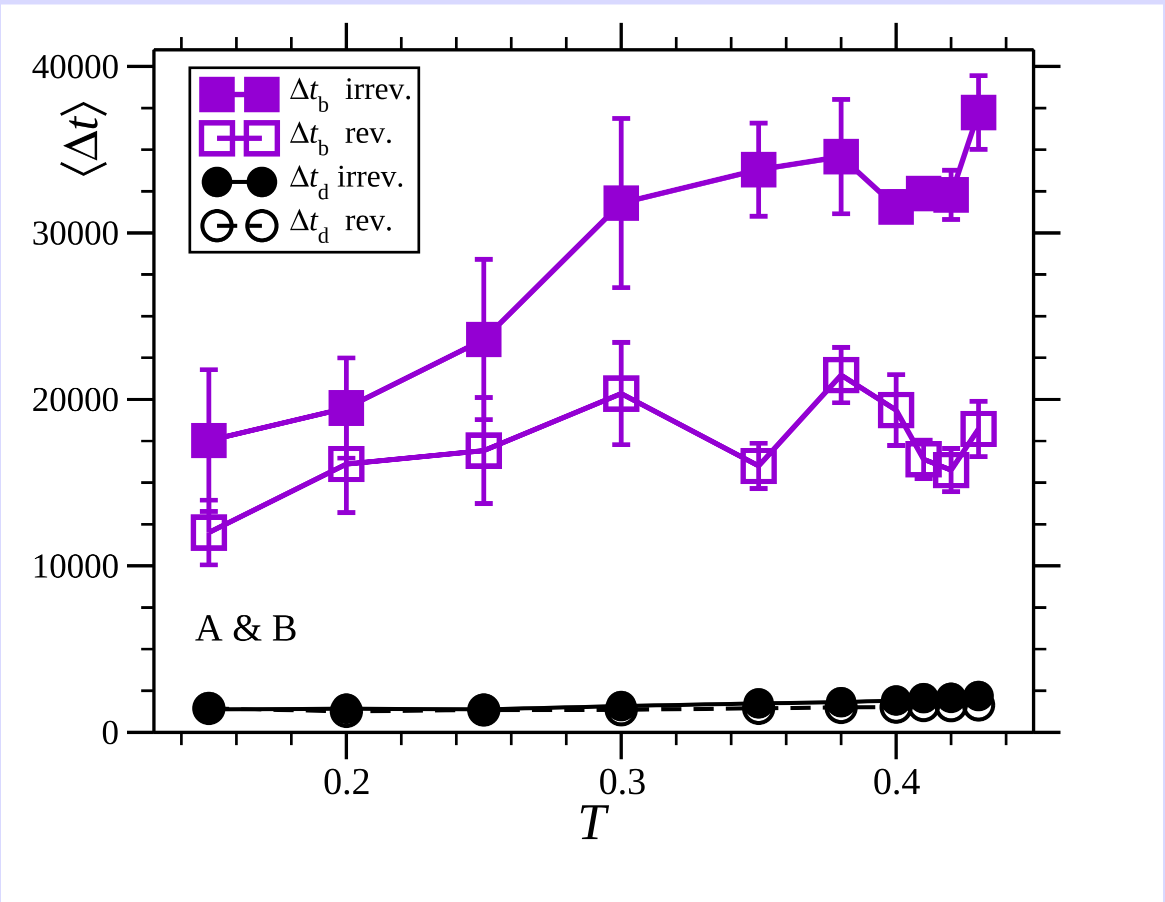
$\Longrightarrow$ irreversible jumps farther

Time Scale



Time Scale





$\Longrightarrow \Delta t_b \gg \Delta t_d$

$\Longrightarrow \Delta t_b$ independent of temperature

(whole simulation 10^5)

Summary: Jump Statistics

At larger temperature relaxation:

- not via Δt_{b} (indep. of T)
- via larger jumpsizes
- via more jumping particles

[J. Chem. Phys. **121**, 4781 (2004)]

Outline

- Jump Statistics
- Correlated Single Particle Jumps
  - ◇ Simultaneously Jumping Particles
  - ◇ Temporally Extended Cluster
- History Dependence
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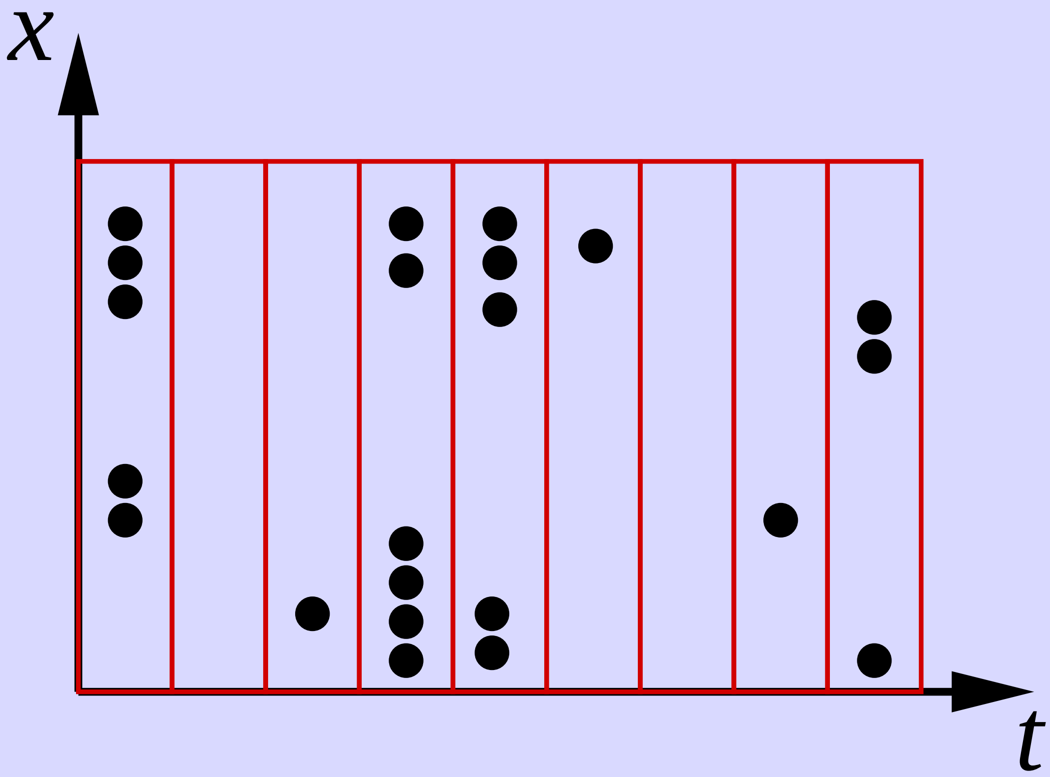
Simultaneously Jumping Particles

Definition: Correlated in Time & Space



Simultaneously Jumping Particles

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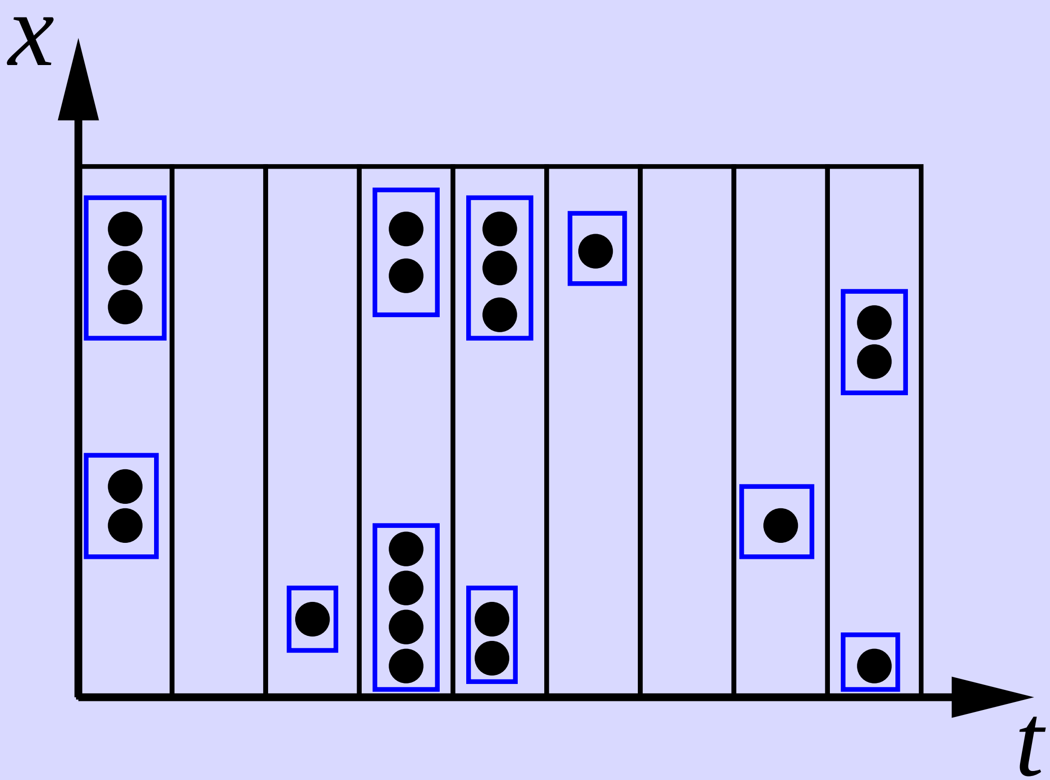
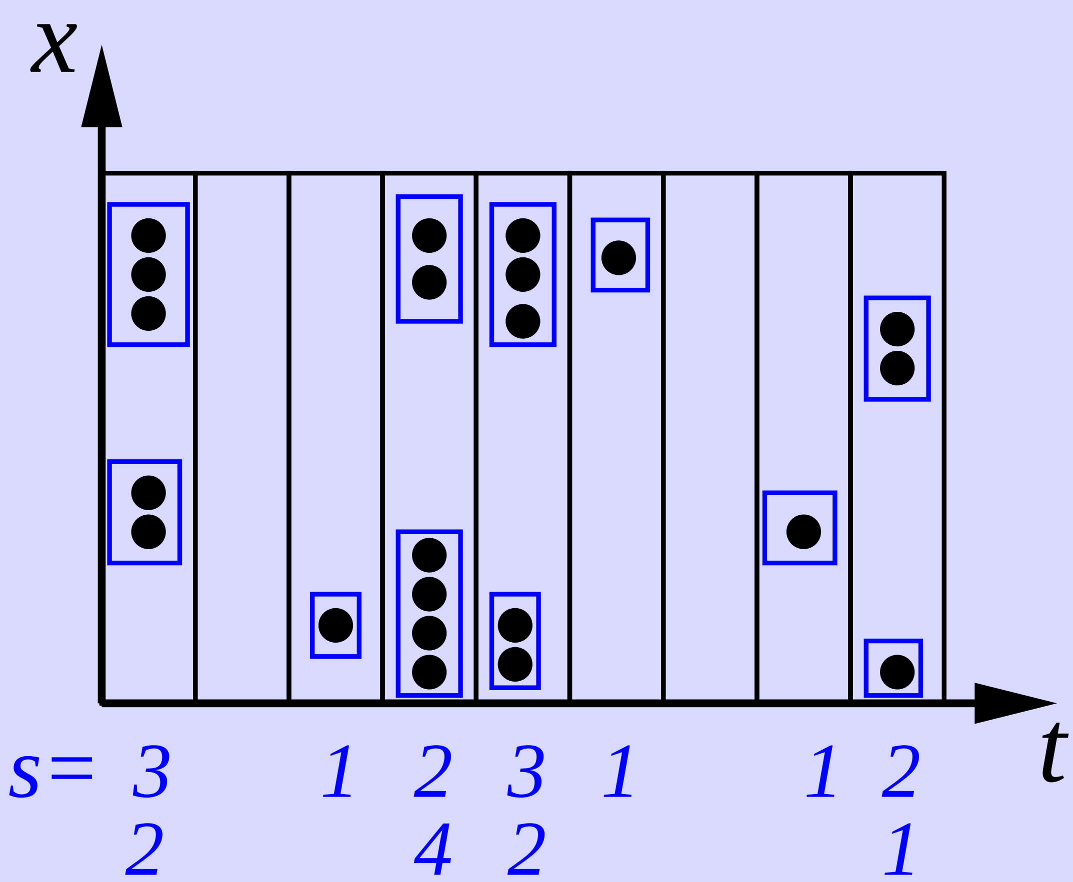


Cluster:

nearest neighbor connections (via $g(r)$)

Simultaneously Jumping Particles

Cluster Size $=$ number of particles in cluster

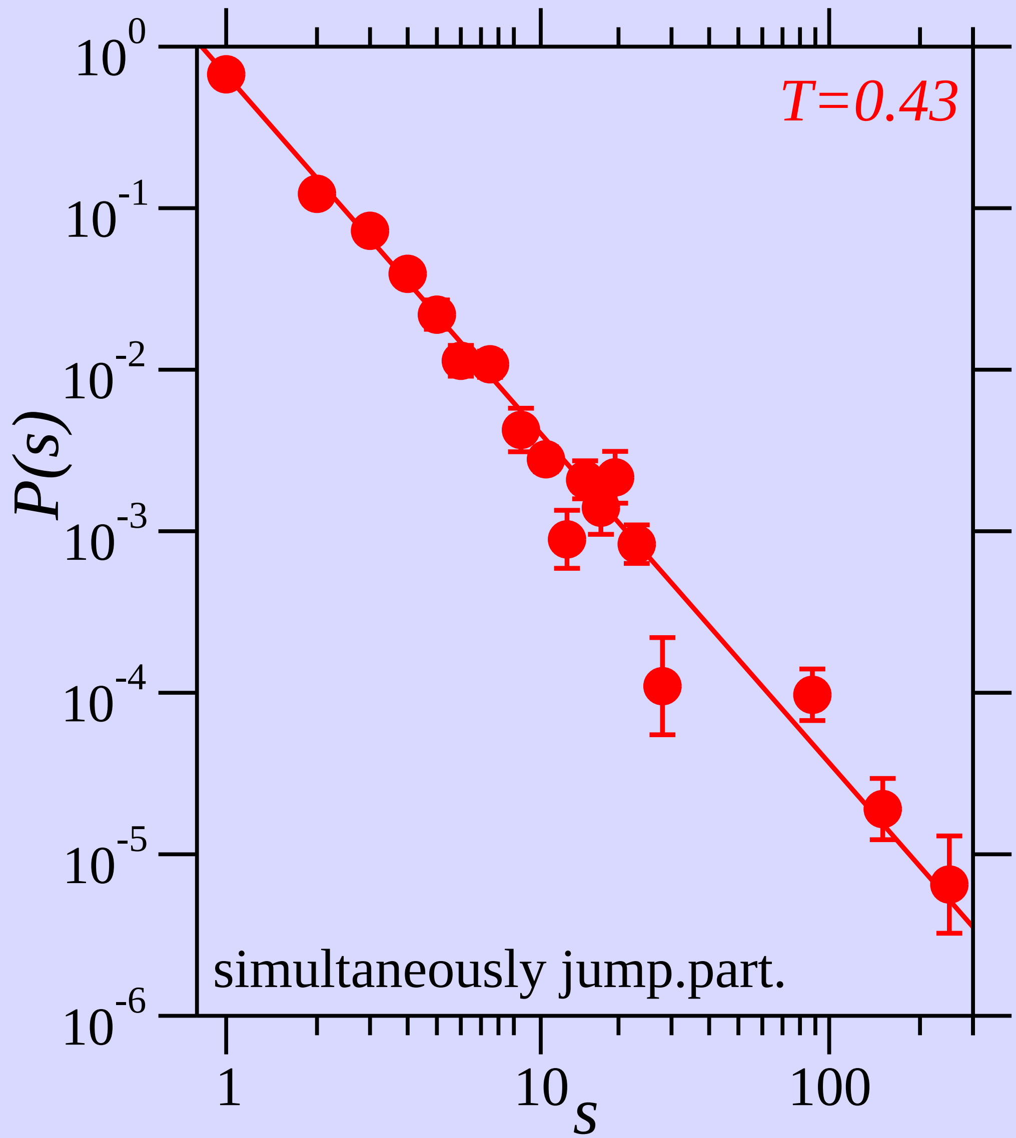


Cluster:

nearest neighbor connections

Cluster Size Distribution

of Simultaneously Jumping Particles



$\Longrightarrow \ln P = a - \tau \ln s$

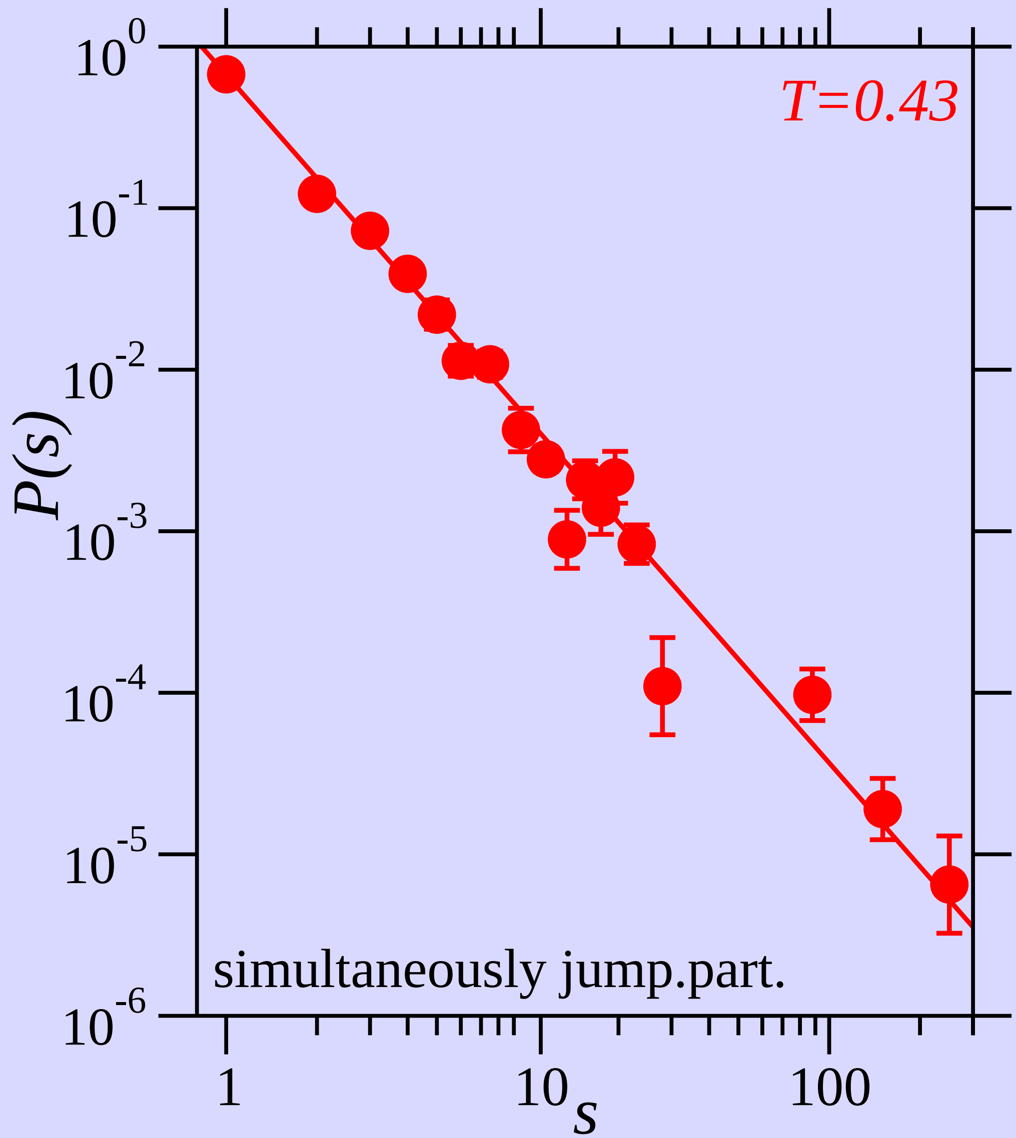
$\Longrightarrow P(s) \sim s^{-\tau}$

$\tau = 1.89 \pm 0.03$

$\Longrightarrow$ critical behavior

Cluster Size Distribution

of Simultaneously Jumping Particles



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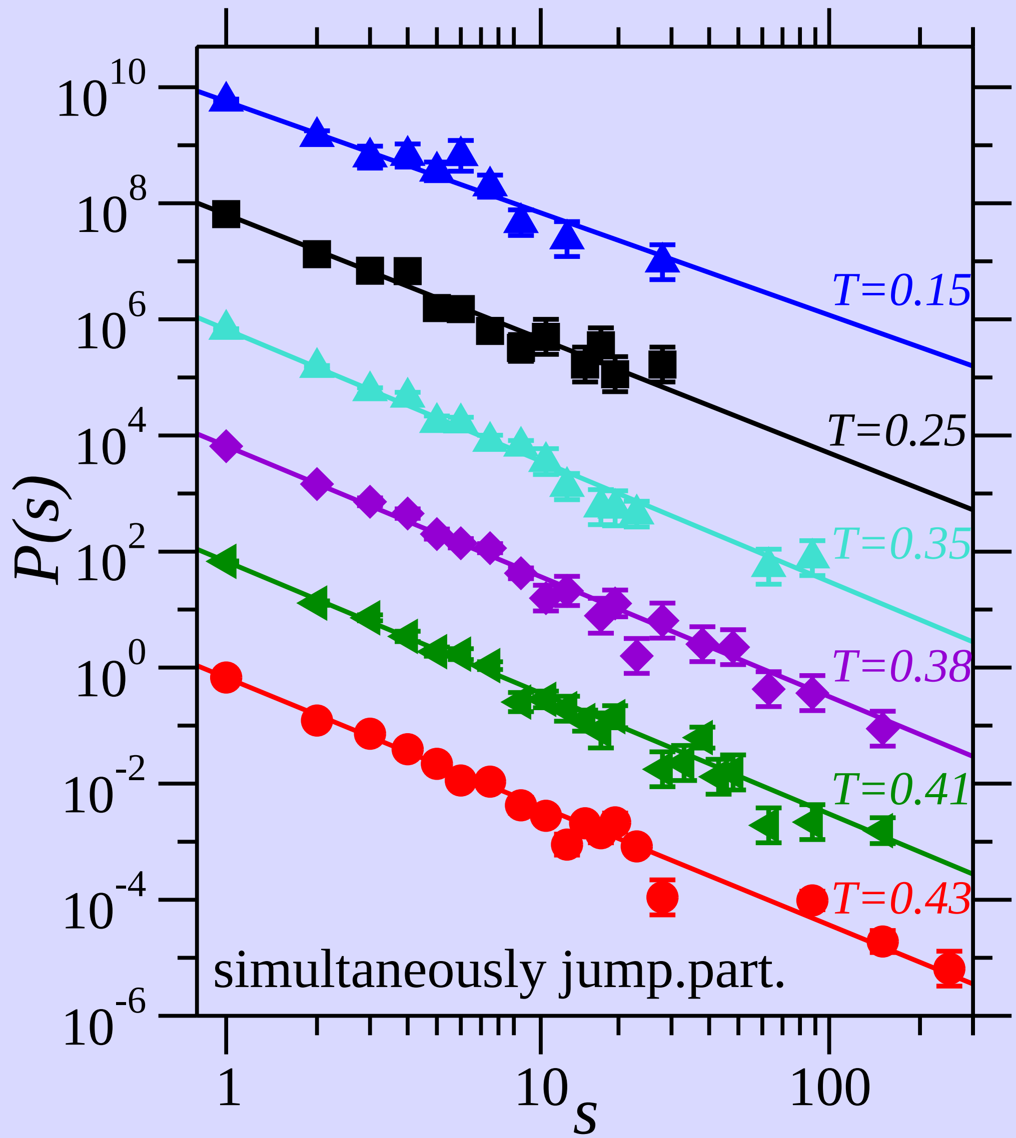
$\Longrightarrow$ critical behavior

$\Longrightarrow$ Percolation?

$\tau_{\mathsf{MF}} = 2.5 \qquad \tau_{3\mathrm{d}} = 2.2$

Cluster Size Distribution

of Simultaneously Jumping Particles



$\Longrightarrow P(s) \sim s^{-\tau}$

percolation?

NO because

$\Longrightarrow$ power law for **all** temperatures

$\Longrightarrow$ **self-organized criticality**

Critical Behavior:

power law $\longleftrightarrow$ scale invariance

Critical Point at Phase Transition:

power law at specific fine tuned external parameter

e.g. percolation: $P(s)=s^{-\tau}$ at $p=p_c$ at all other p no power law

e.g. jumping particle clusters: $P(s)=s^{-\tau}$ (would be) at $T=T_c$ only

Self−Organized Criticality:

power law for whole range of external parameter

here jumping particle clusters: $P(s)=s^{-\tau}$ for all $T=0.15-0.43$

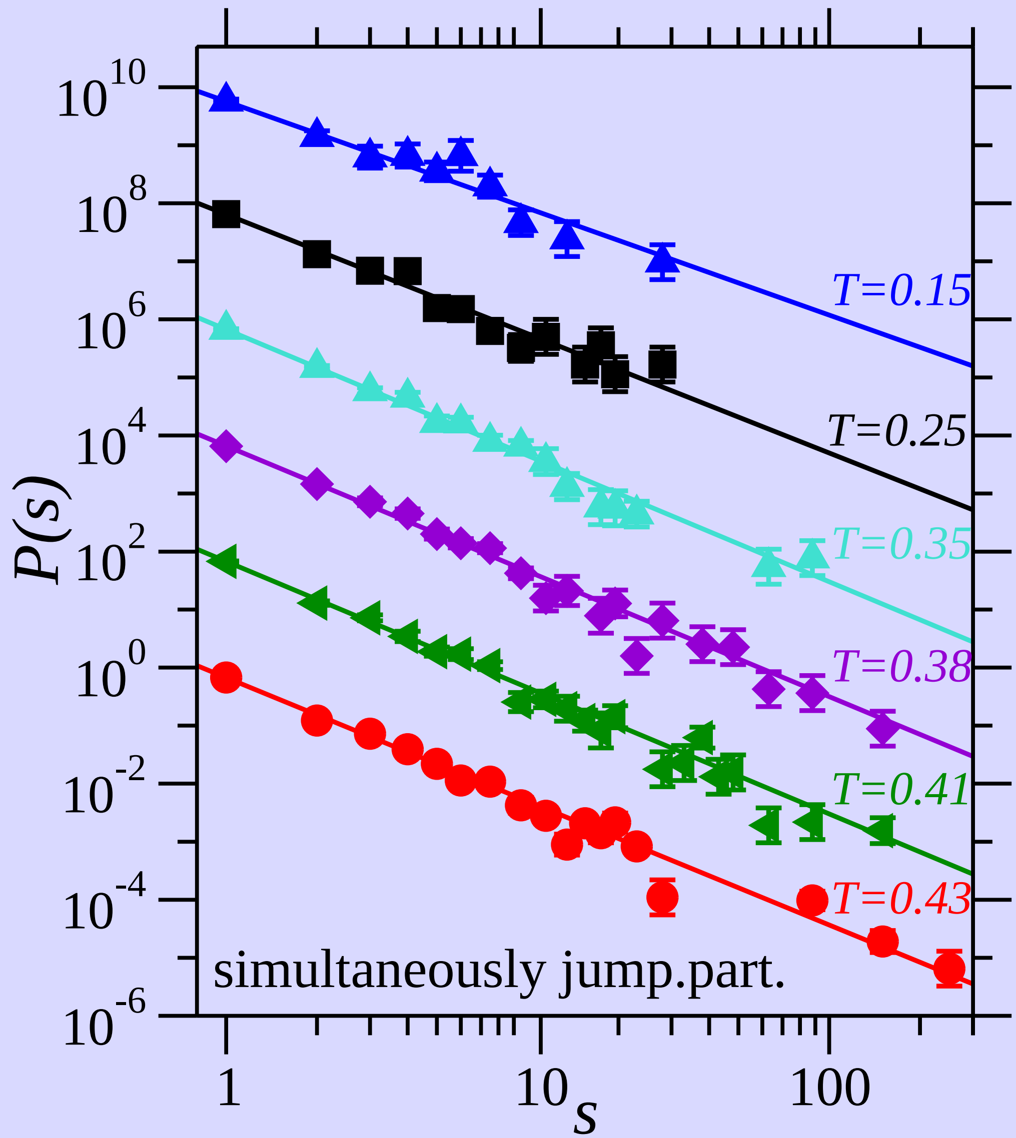
Other Examples:

- sandpile avalanches
- forest fire
- financial market
- earth quakes

[P. Bak, C. Tang, and K. Wiesenfeld, PRL 59, 381 (1987)]

Cluster Size Distribution

of Simultaneously Jumping Particles



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$\tau(T)$

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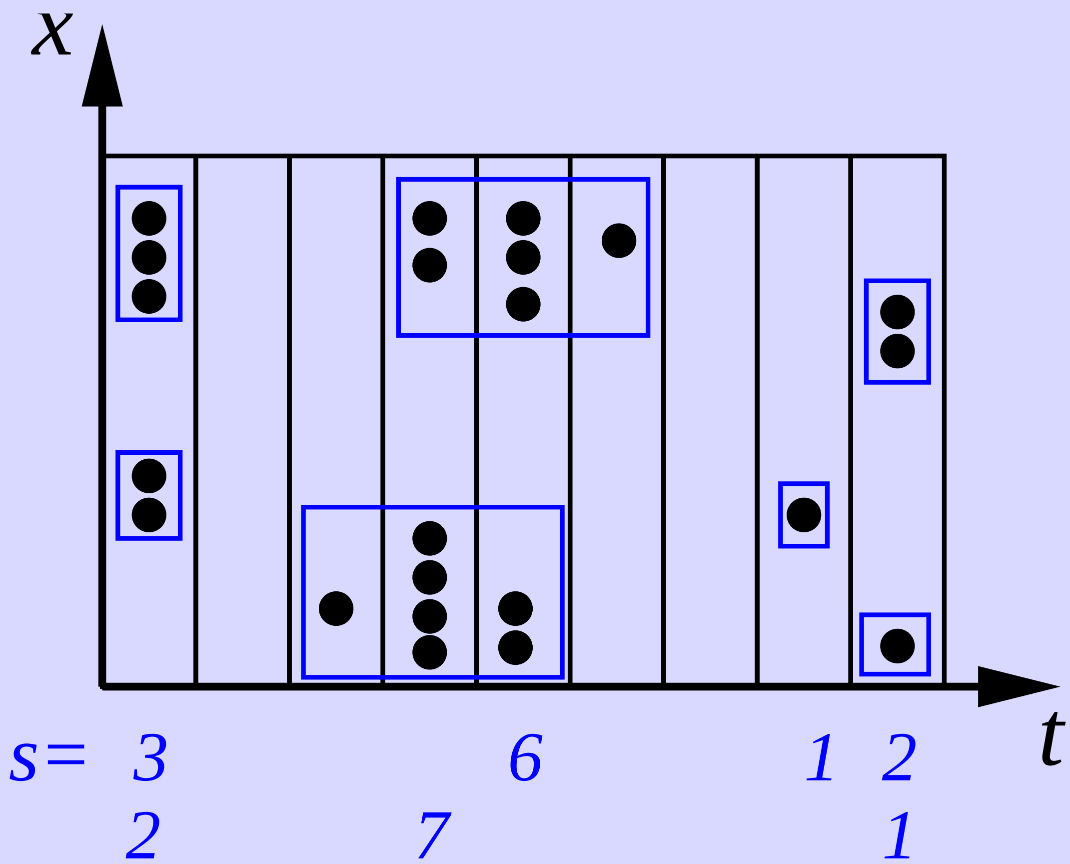
Temporally Extended Cluster



Definition:

cluster of events $(\mathbf{r}_i, t_i)$

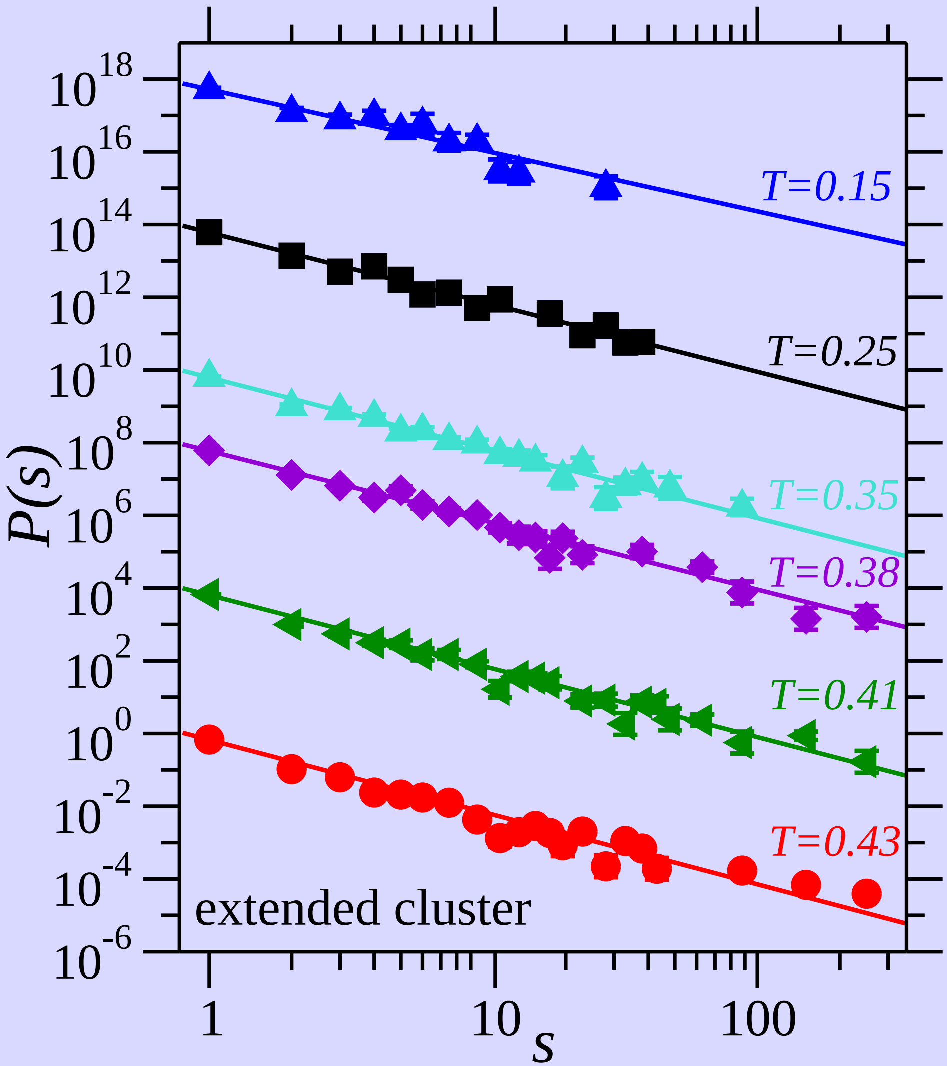
connected if:

$$\Delta r < r_{\text{cutoff}} \quad \text{and}$$

$$\Delta t < t_{\text{cutoff}}$$

Cluster Size Distribution

of Temporally Extended Clusters

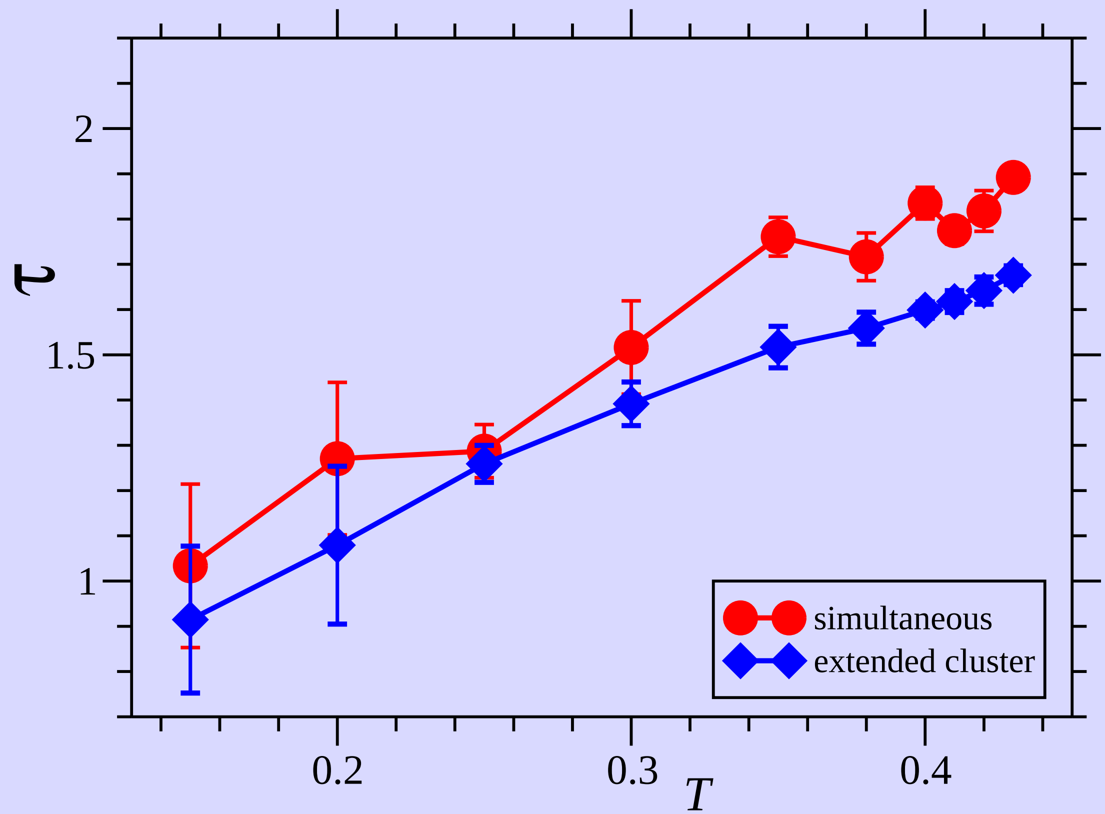


$\Longrightarrow P(s) \sim s^{-\tau}$

$\Longrightarrow$ for all temperatures (self-organized crit.)

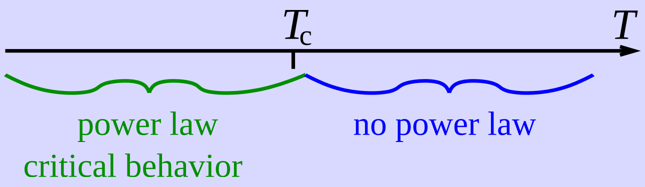
Exponent $\tau(T)$

$$P(s) \sim s^{-\tau}$$



slightly above T_c $\tau \approx 1.86$

[Donati et al. 1999]



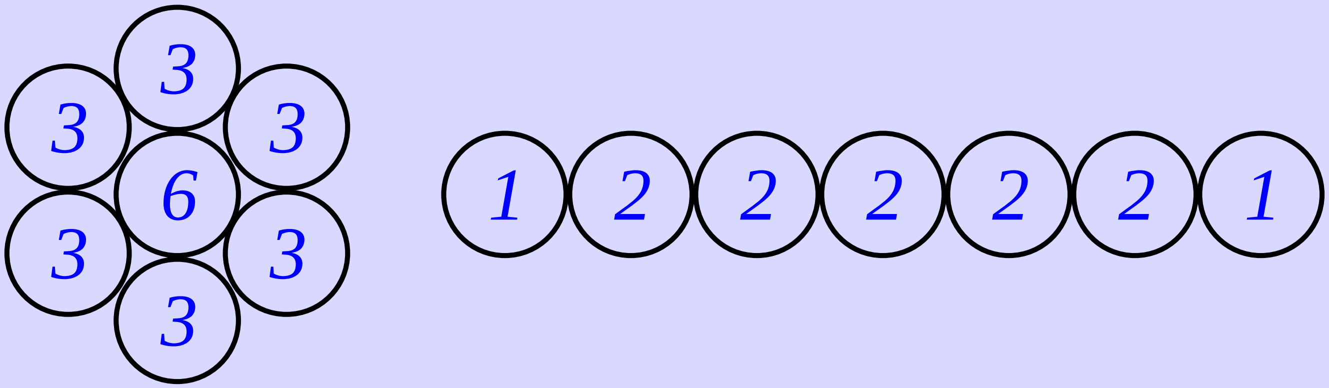
P simult.

Shape of Clusters

z = number of nearest neighbors within cluster

s = number of particles (cluster size)

$\langle z \rangle$ = average of z over particles $1, \ldots s$



$s=7$ $\langle z \rangle=3.4$

$s=7$ $\langle z \rangle=1.7$

Shape of Clusters

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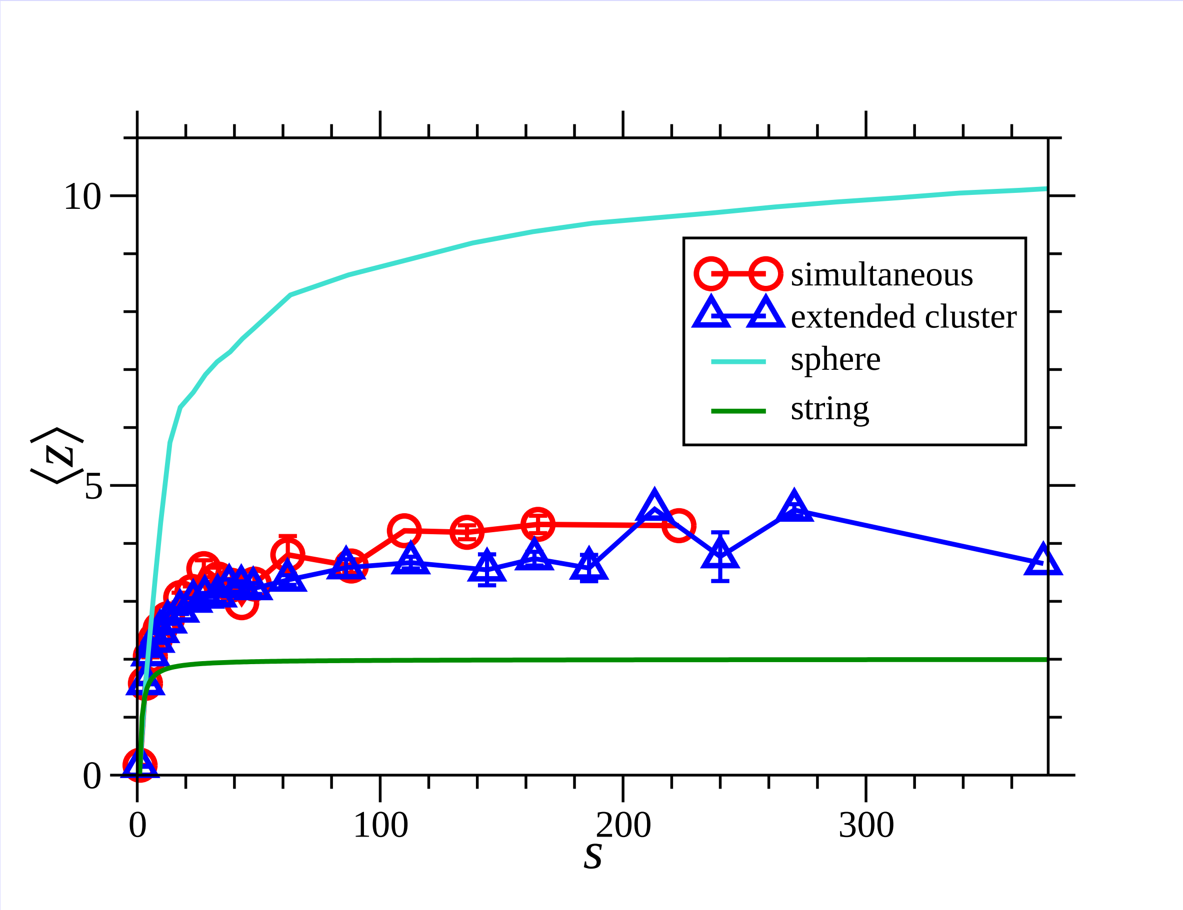

$\Longrightarrow$ string-like clusters

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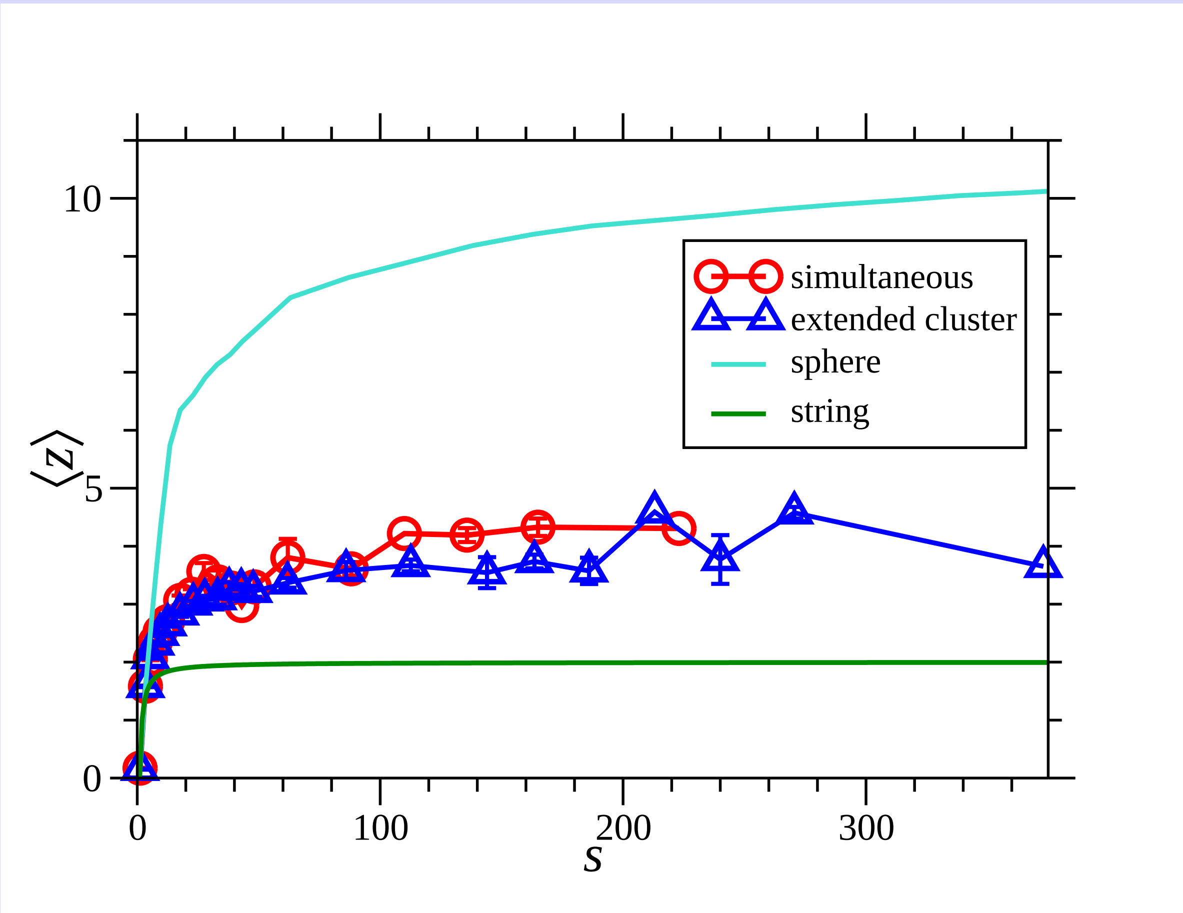
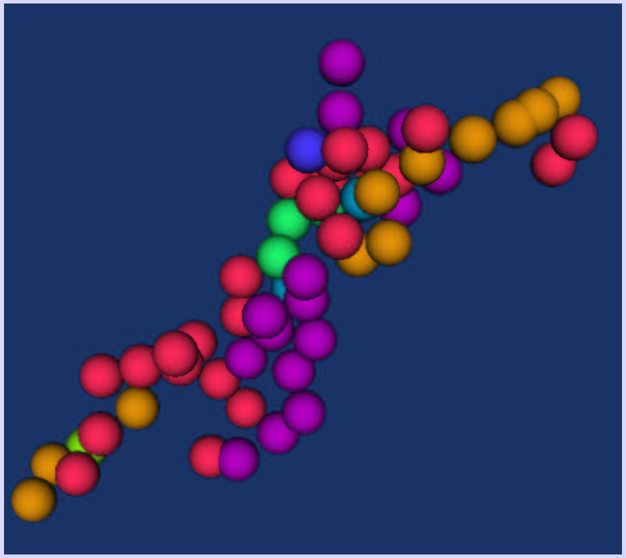


$\Longrightarrow$ string-like clusters



same color = same time

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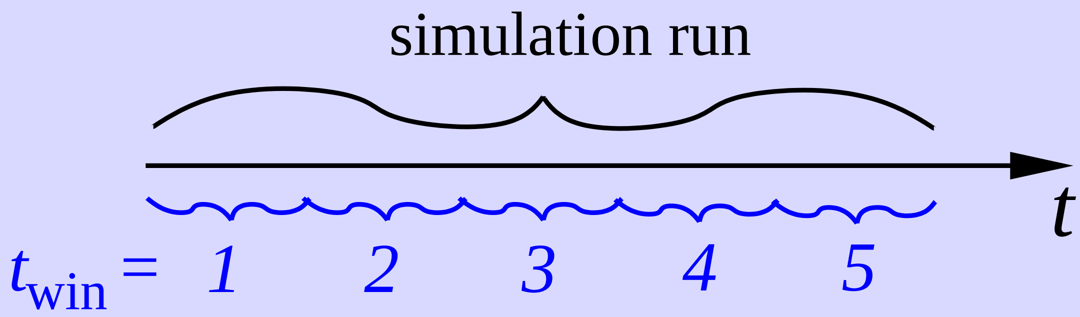
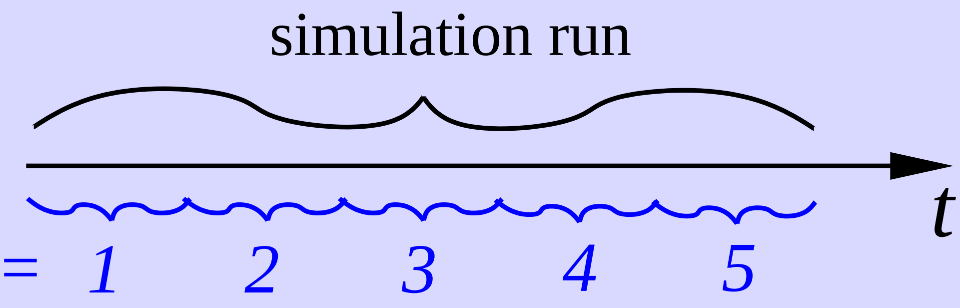
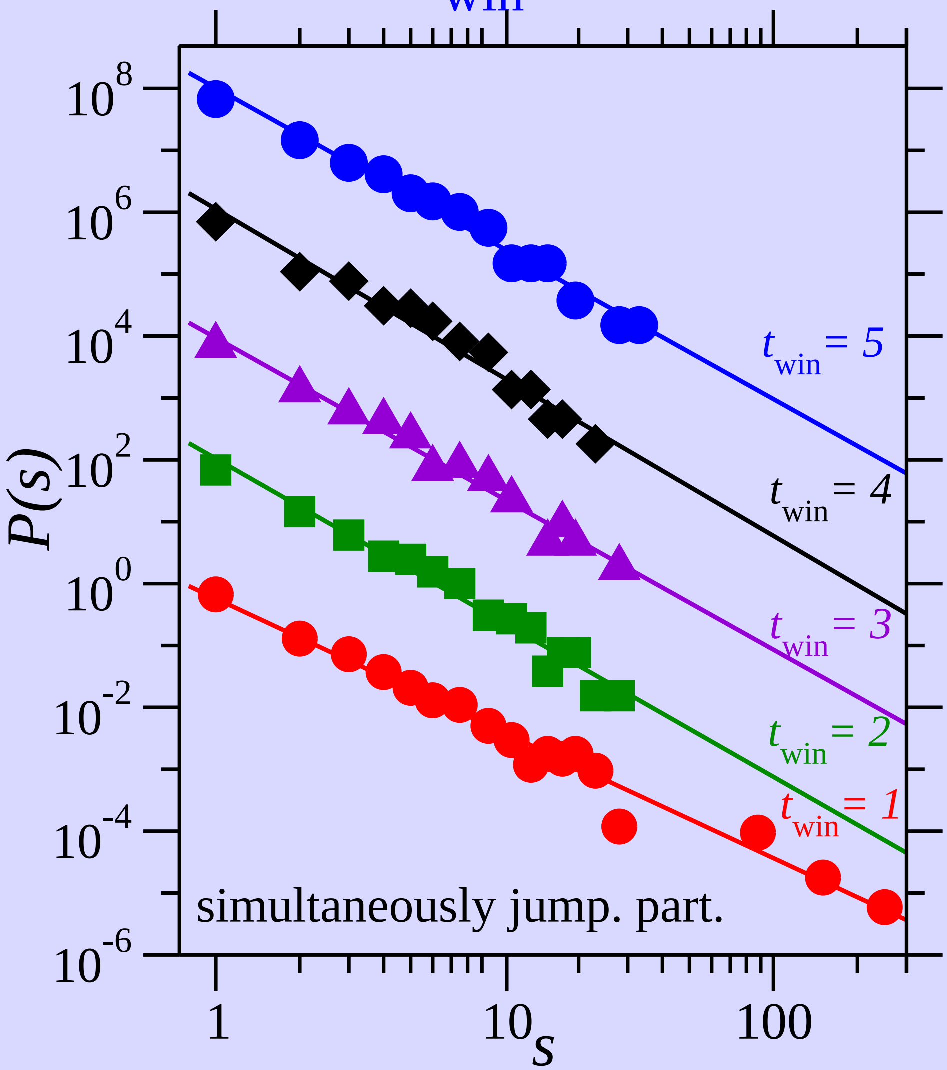
History Dependence



History Dependence





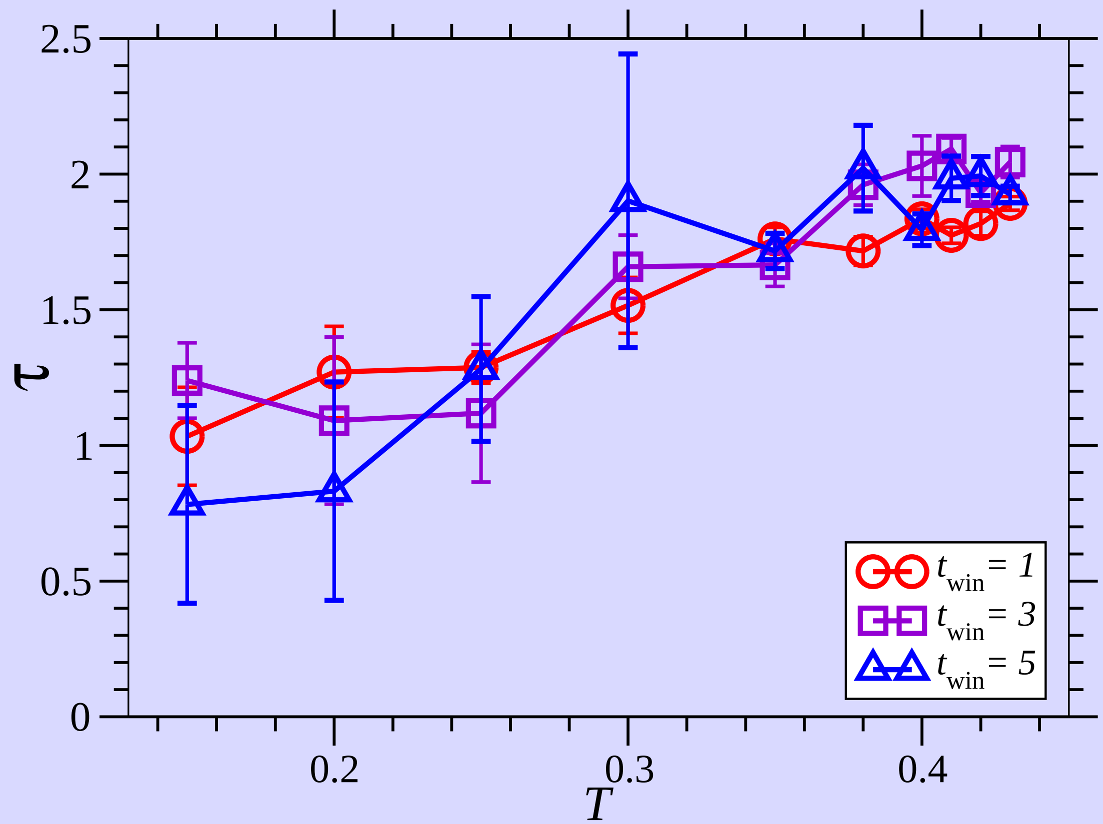
Power Law

(Simultan. Jump. Part.)

$\Longrightarrow$ aging independent

Exponent $\tau(T)$

$$P(s) \sim s^{-\tau}$$



$\Longrightarrow$ aging independent

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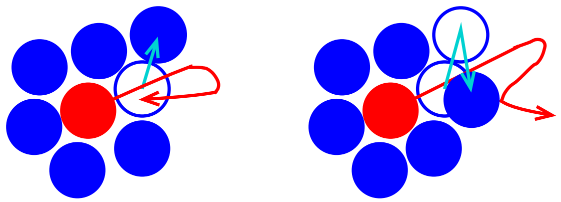
Summary: Jump Statistics

reversible and irreversible jumps:



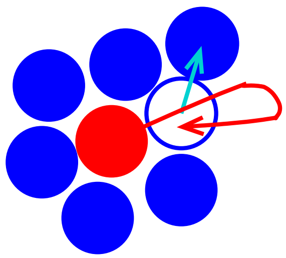
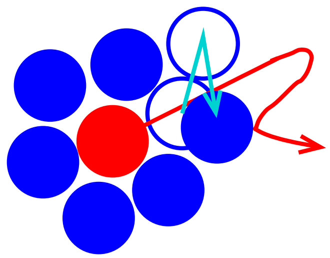
reversible jump irreversible jump

At larger temperature relaxation:

- via more jumping particles
- via larger jumpsizes
- not via Δt_{b} (indep. of T)

Summary: Jump Statistics

reversible and irreversible jumps:

reversible jump irreversible jump

At larger temperature relaxation:

- via more jumping particles history dependent
- via larger jumpsizes history independent
- not via Δt_{b} (indep. of T) history independent

Summary: Correlated Single Particle Jumps

simultaneously jump. part. & extended clusters

- jumps are correlated spatially and temporally
- Distribution of Cluster Size: $P(s) \sim s^{-\tau}$
  - ◇ aging independent
  - ◇ for all temp. $\longrightarrow$ **self-organized criticality** (critical behavior gets frozen in)
- string-like clusters

[Europhys. Lett. **76**, 1130 (2006)]

Future/Present

- SiO_2
  (R. A. Bjorkquist & J. A. Roman & J. Horbach)
- granular media
  (T. Aspelmeier & A. Zippelius)

Acknowledgments

A. Zippelius, K. Binder, E. A. Baker, J. Horbach

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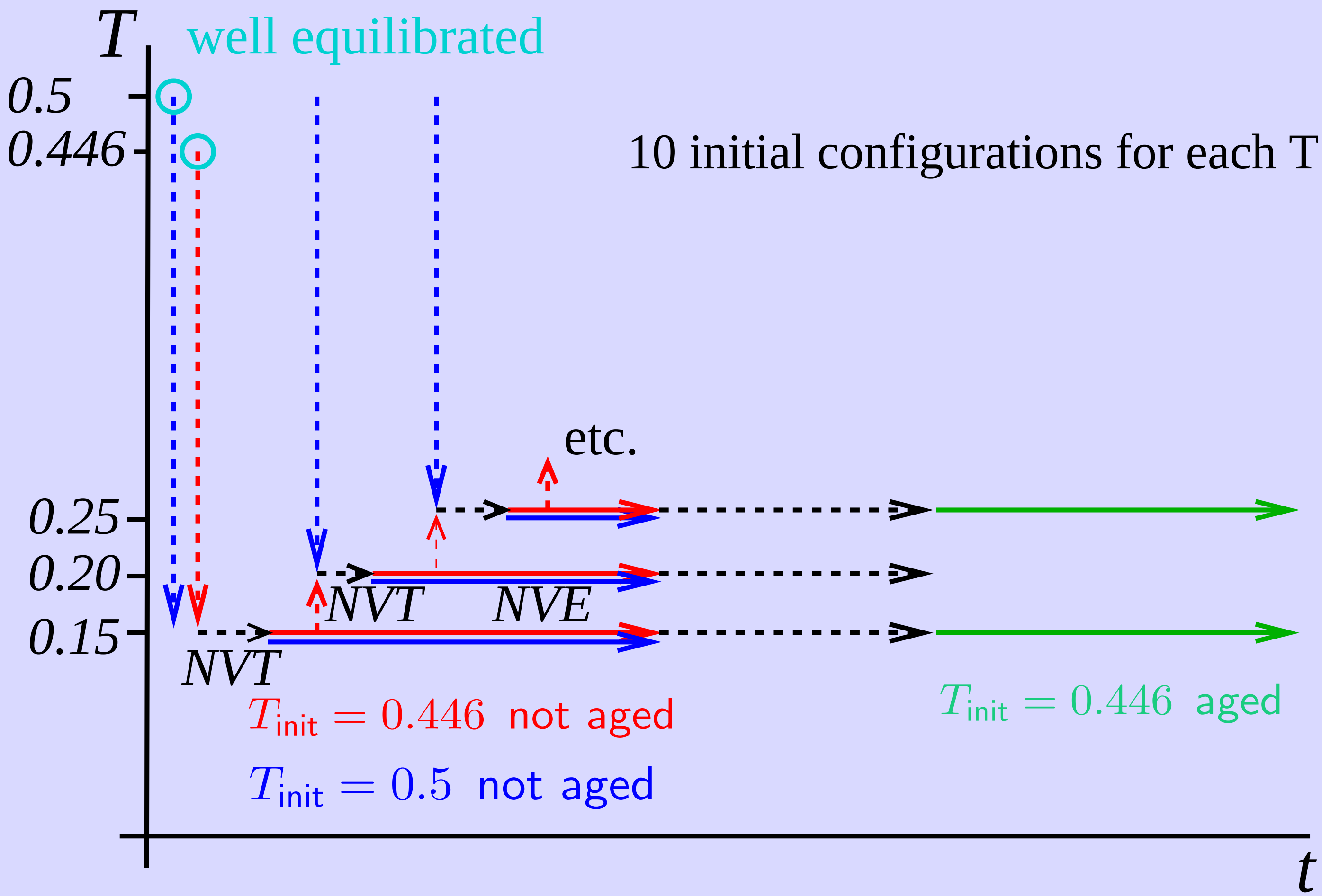
SFB 262 and DFG Grant No. Zi 209/6-1

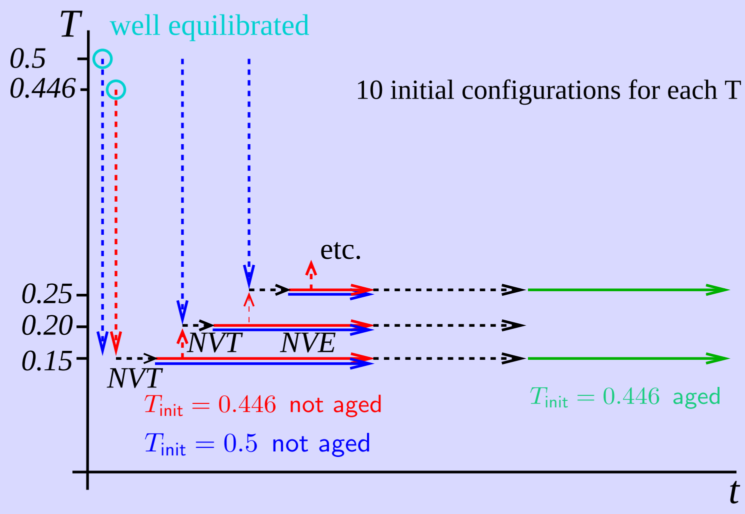
Time Scales

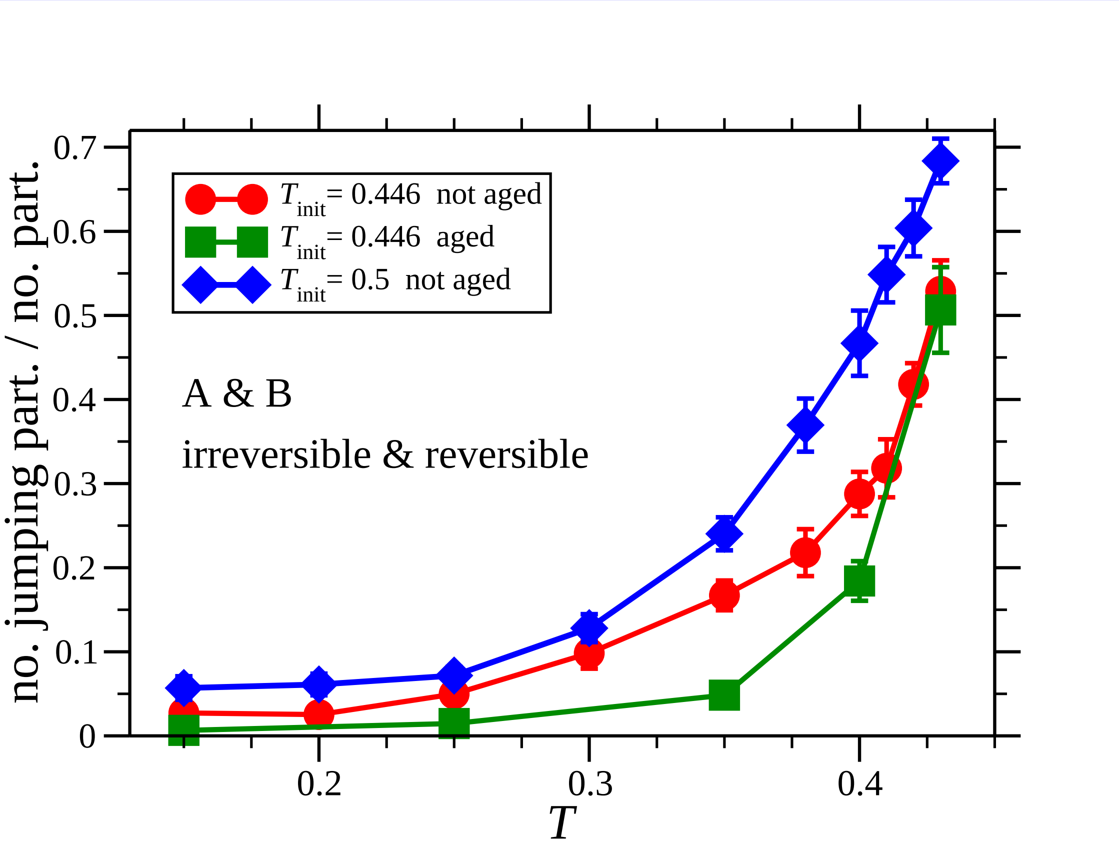
- one MD step: 0.02 time units, Ar: $3 \cdot 10^{-13}$s$\cdot 0.02 = 6$fs
- one oscillation: 100 MD steps, 0.6 ps
- time a jump takes: 200 MD steps, 1.2 ps
- time resolution (time bin): 40000 MD steps, 240 ps
- time betw. successive jumps Δt_{b}: $1.5 \cdot 10^6$ MD steps, 9 ns
- whole simulation run: $5 \cdot 10^6$ MD steps, 30 ns

History of Production Runs



History Dependence

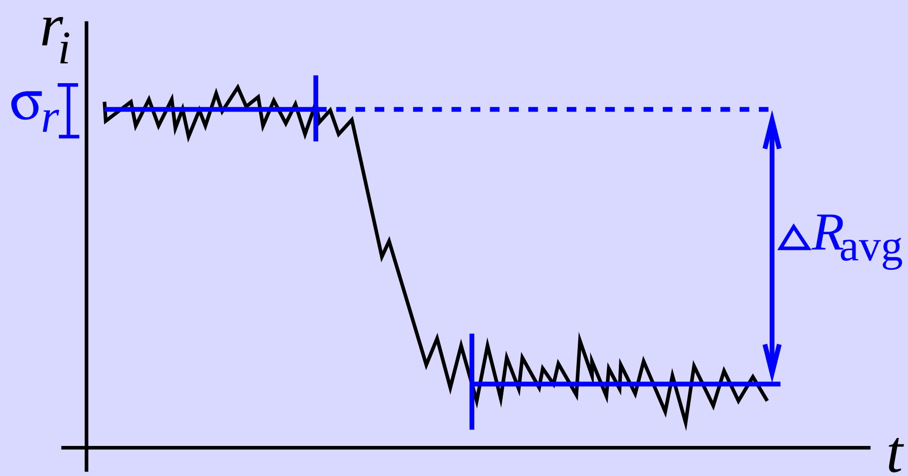


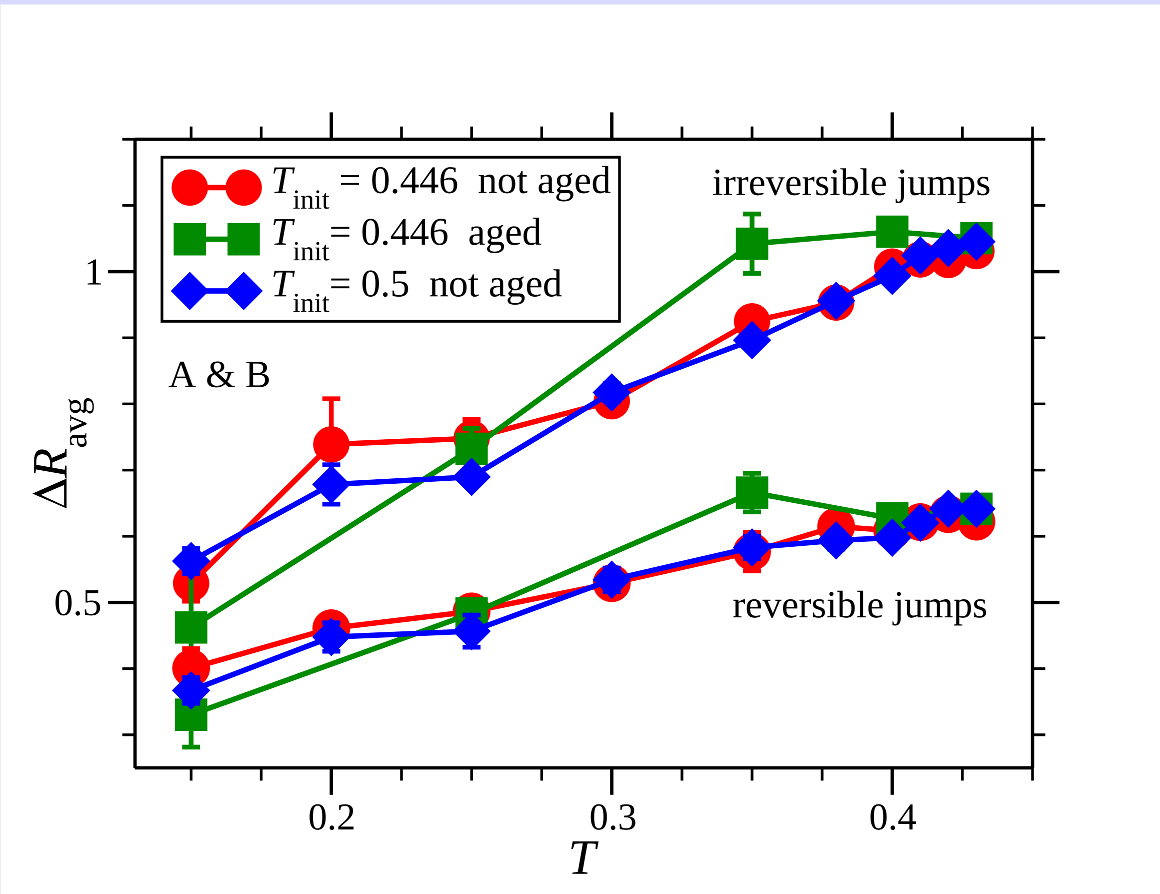


Number of Jump. Part.

$\Longrightarrow$ history dependent

History Dependence

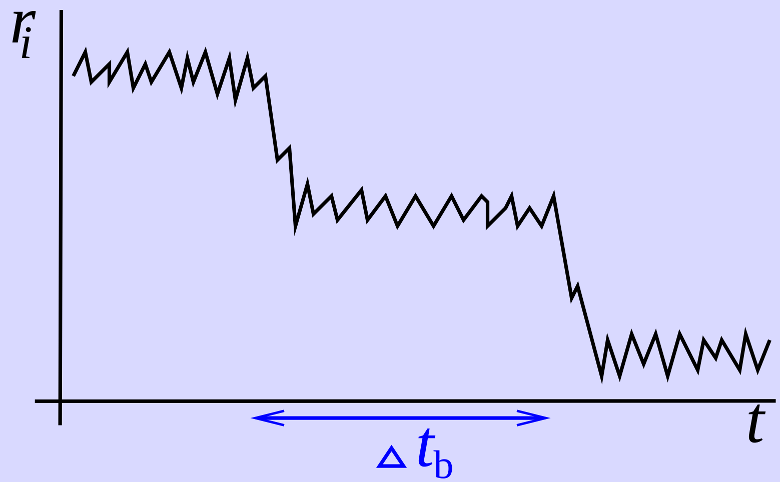


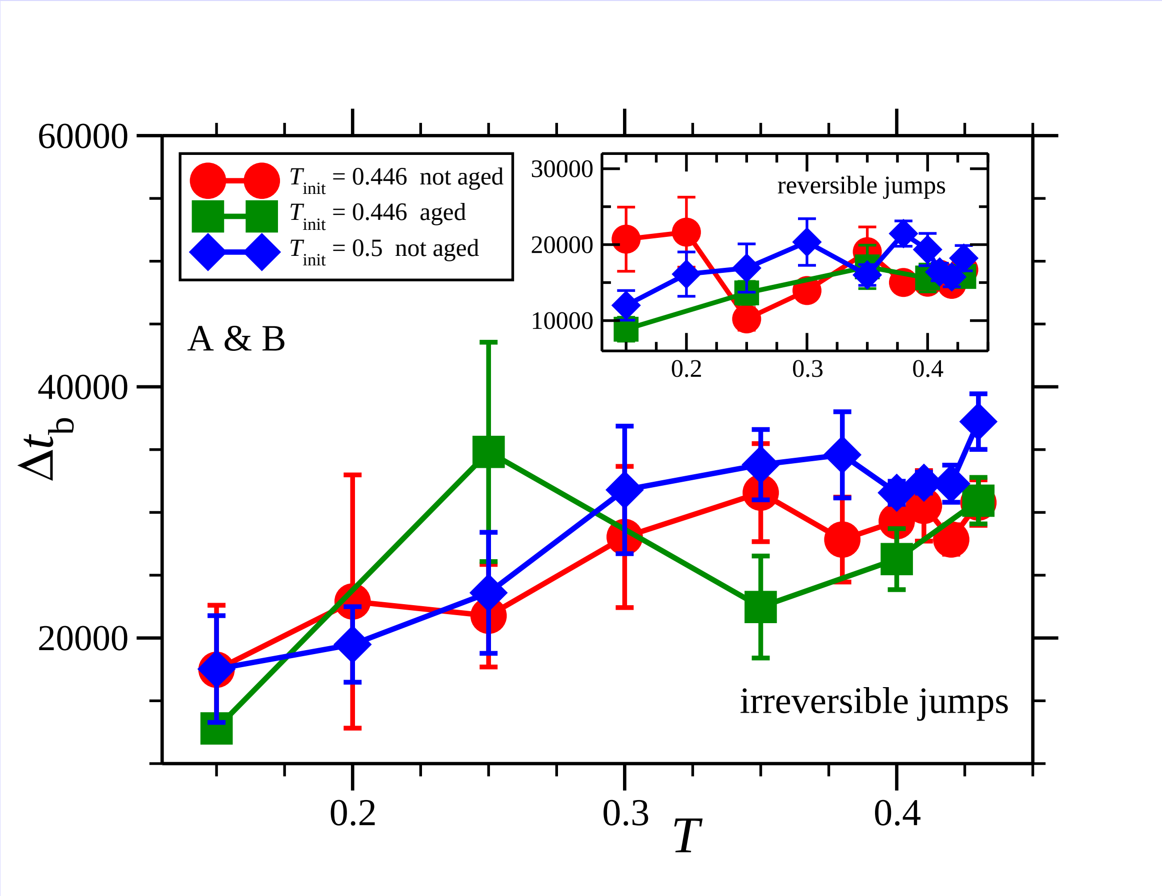


Jump Size

$\Longrightarrow$ history independent

History Dependence

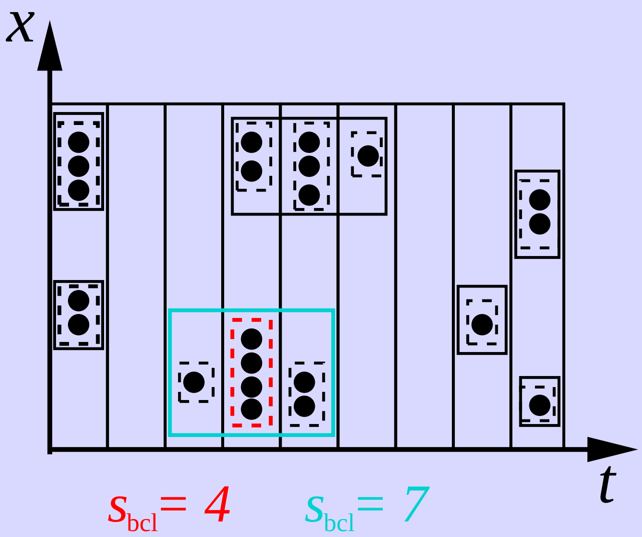




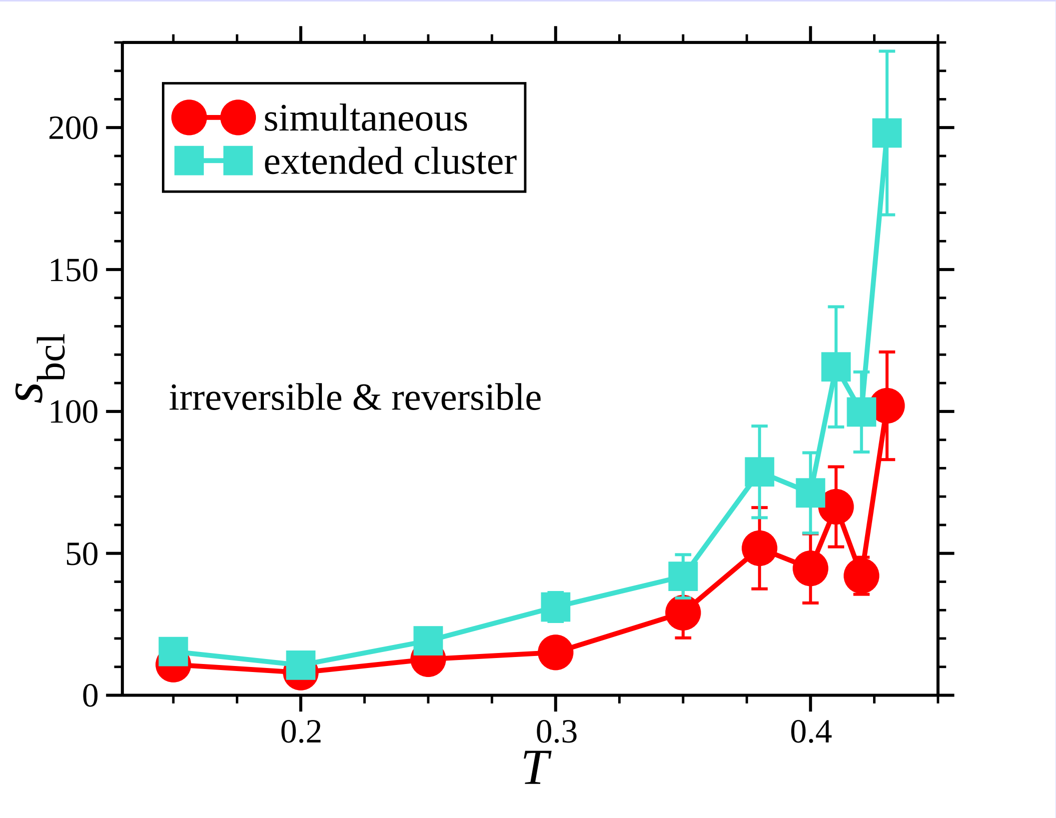
Time Between Jumps

$\Longrightarrow$ history independent

Most Cooperative Processes



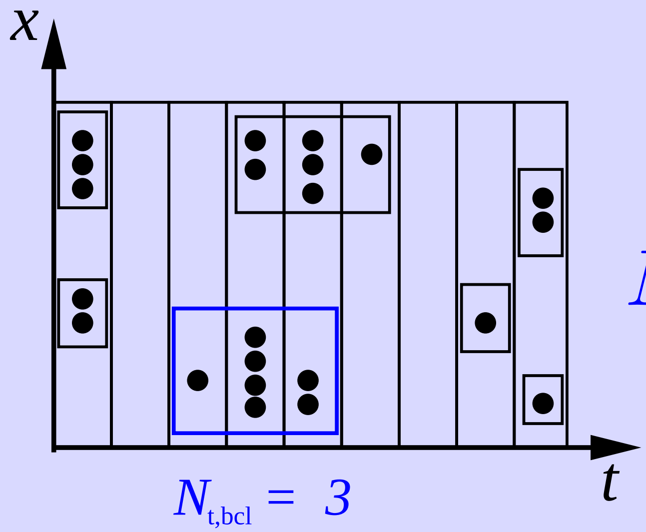
s_{bcl} = largest cluster size



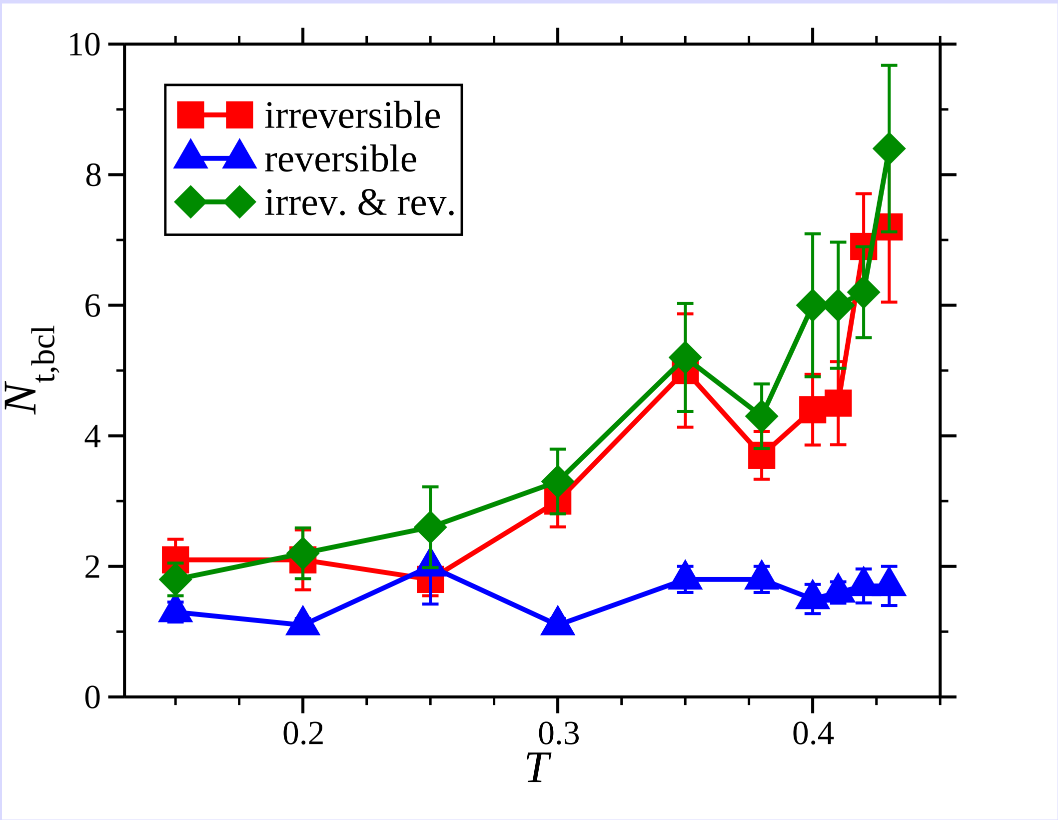
$\Longrightarrow$ highly correlated single particle jumps

- many particles

Most Cooperative Processes



$N_{t,bcl}$ = no. of time bins of largest cluster

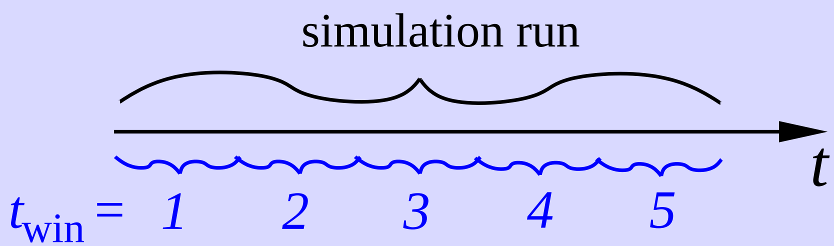


$\Longrightarrow$ highly correlated single particle jumps

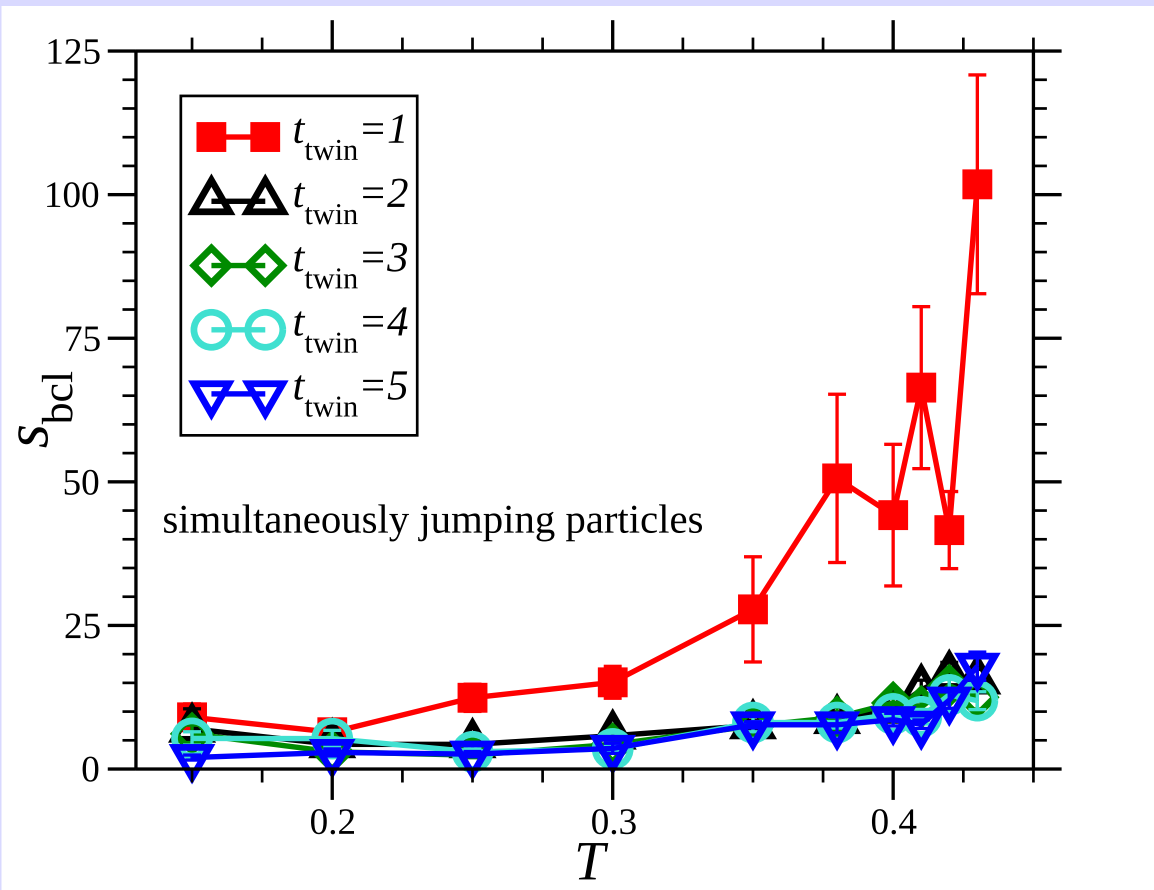
- many particles
- many time bins

(maximum = 125)

History Dependence



s_{bcl} = largest cluster size

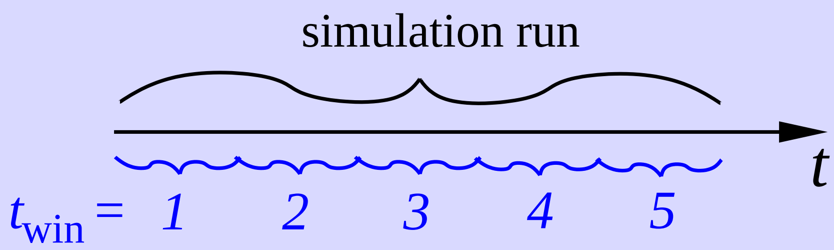


$\Longrightarrow$ aging dependent

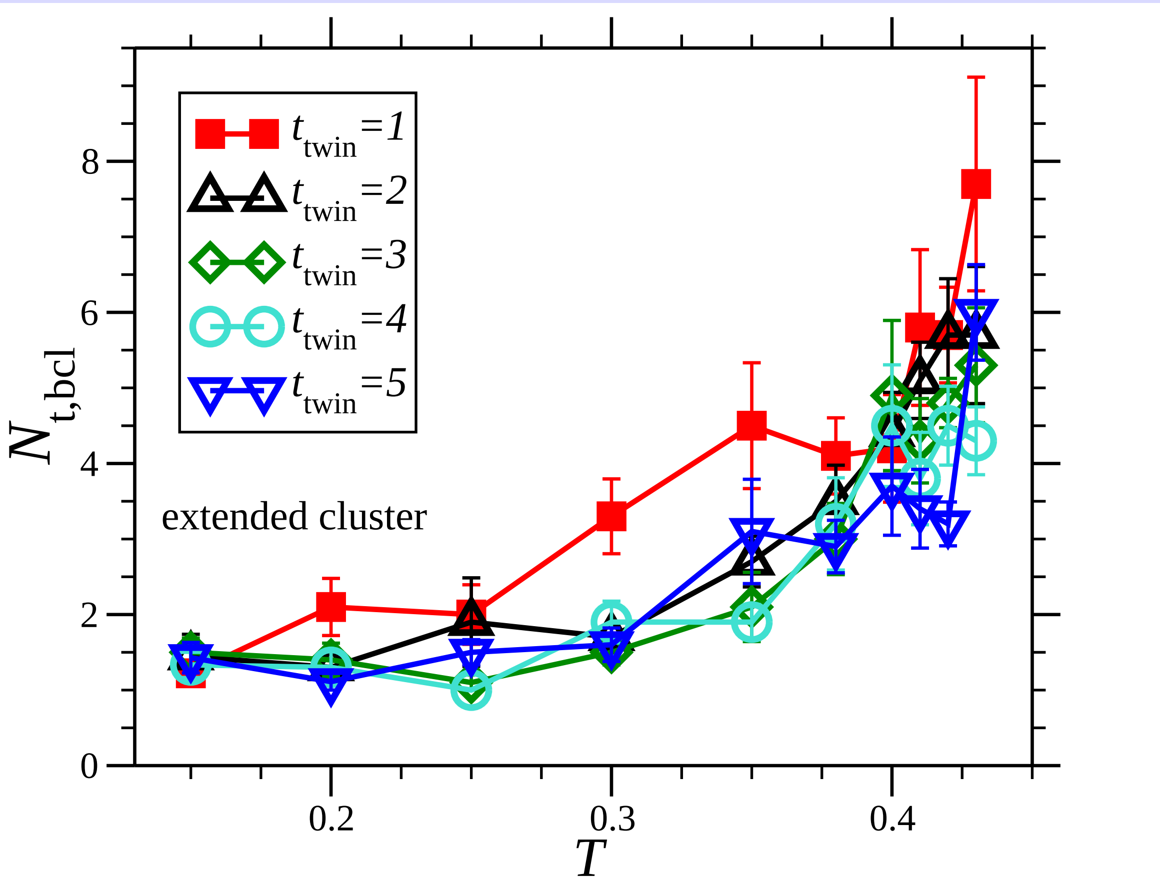
- 1st t-window:
  highly cooperative
- 2nd - 5th t-window:
  same, cooperative

s_{bcl} extended cluster

History Dependence



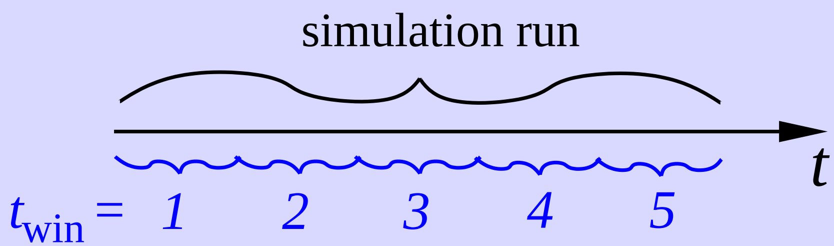
$N_{\mathrm{t,bcl}}$ = no. of t-bins of largest cluster

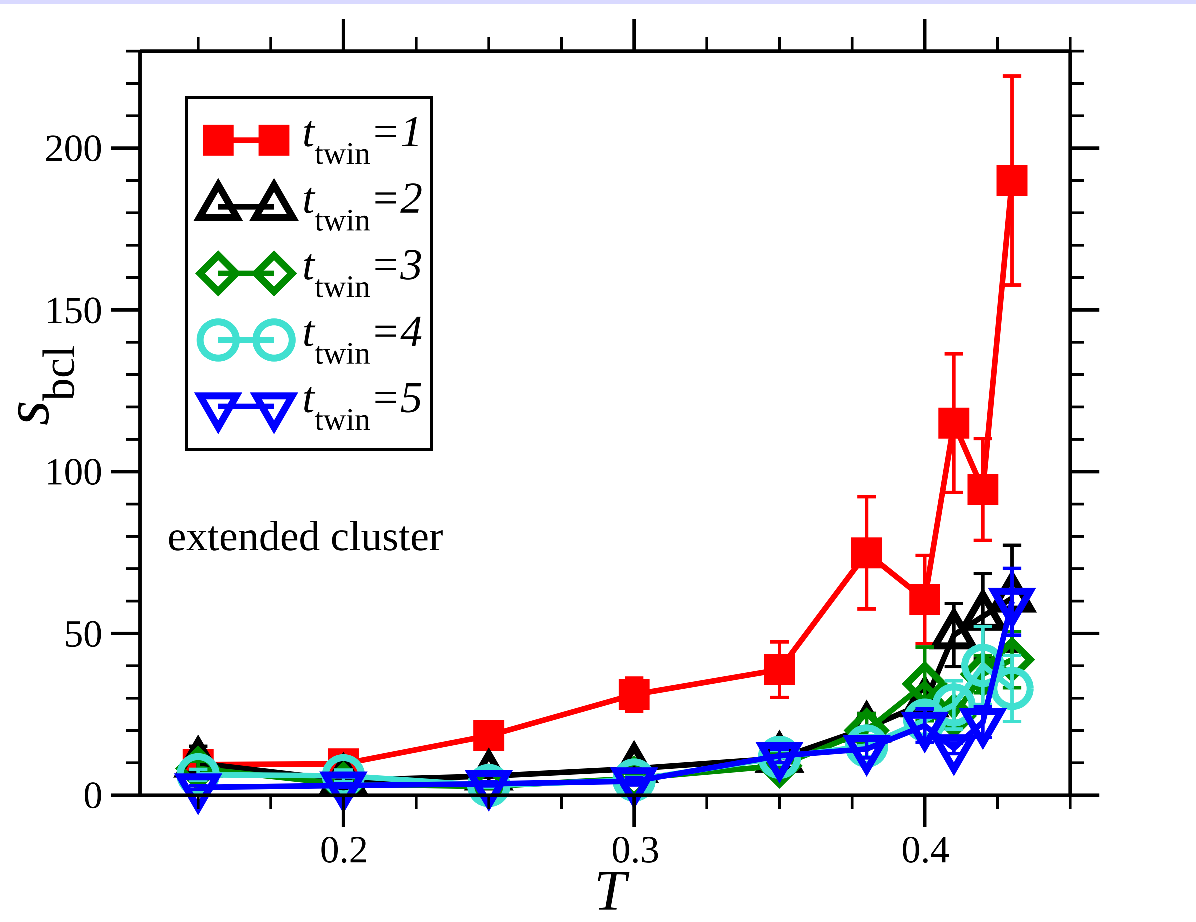


$\Longrightarrow$ less aging dependent

$\Longrightarrow$ highly cooperative

History Dependence





$\Longrightarrow$ aging dependent

- 1st t-window:
  highly cooperative
- 2nd - 5th t-window:
  same, cooperative

s_{bcl} simult. jump.

Molecular Dynamics Simulation

Initialize:

$\vec{x}_i(t_0),\ \vec{v}_i(t_0),\ \vec{a}_i(t_0)$

particles $i=1,\dots,N$

three dimensions

↓

$\vec{x}_i(t_0+\Delta t),\ \vec{v}_i(t_0+\Delta t),\ \vec{a}_i(t_0+\Delta t)$

↓

$\vec{x}_i(t_0+2\Delta t),\ \vec{v}_i(t_0+2\Delta t),\ \vec{a}_i(t_0+2\Delta t)$

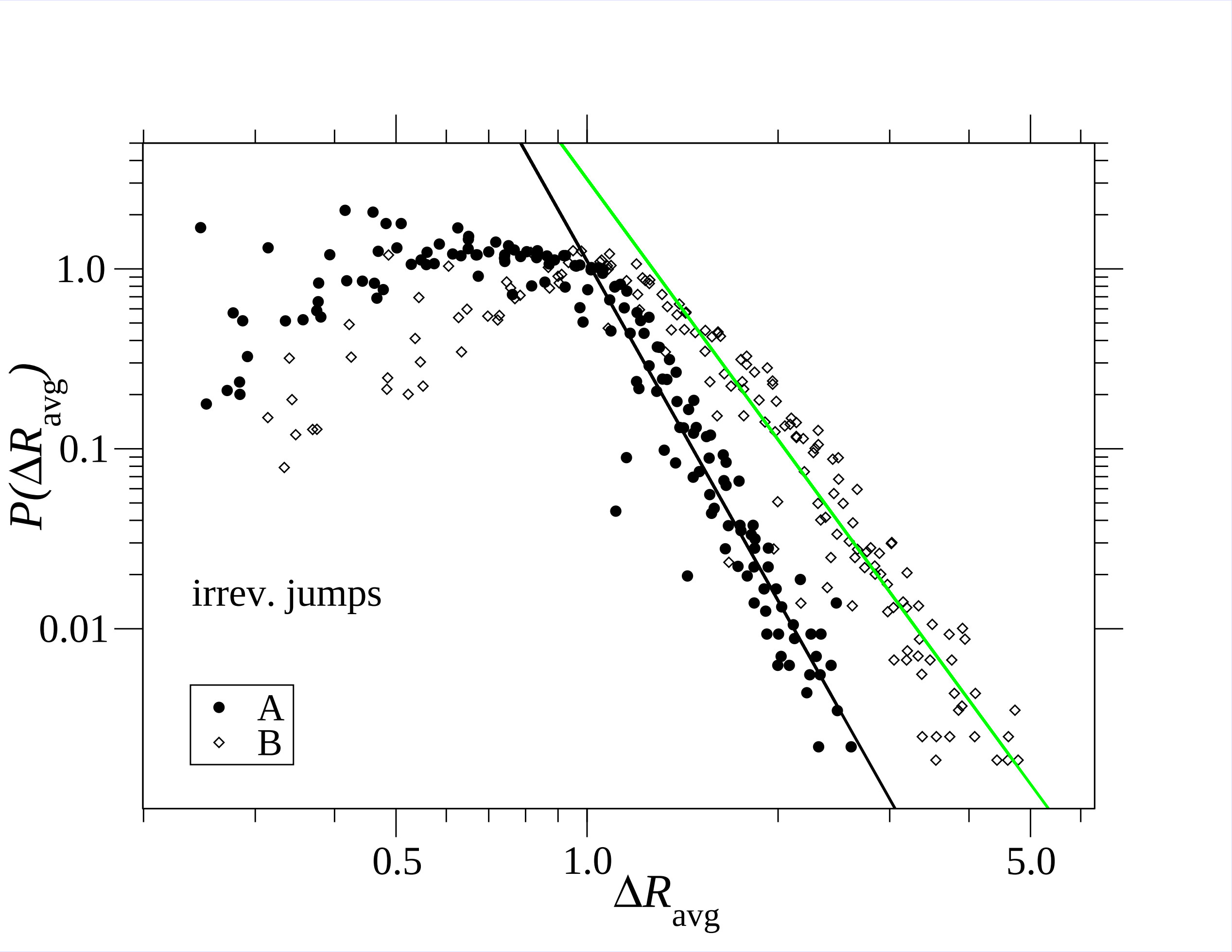
↓ etc.

↓ = Iteration Step:

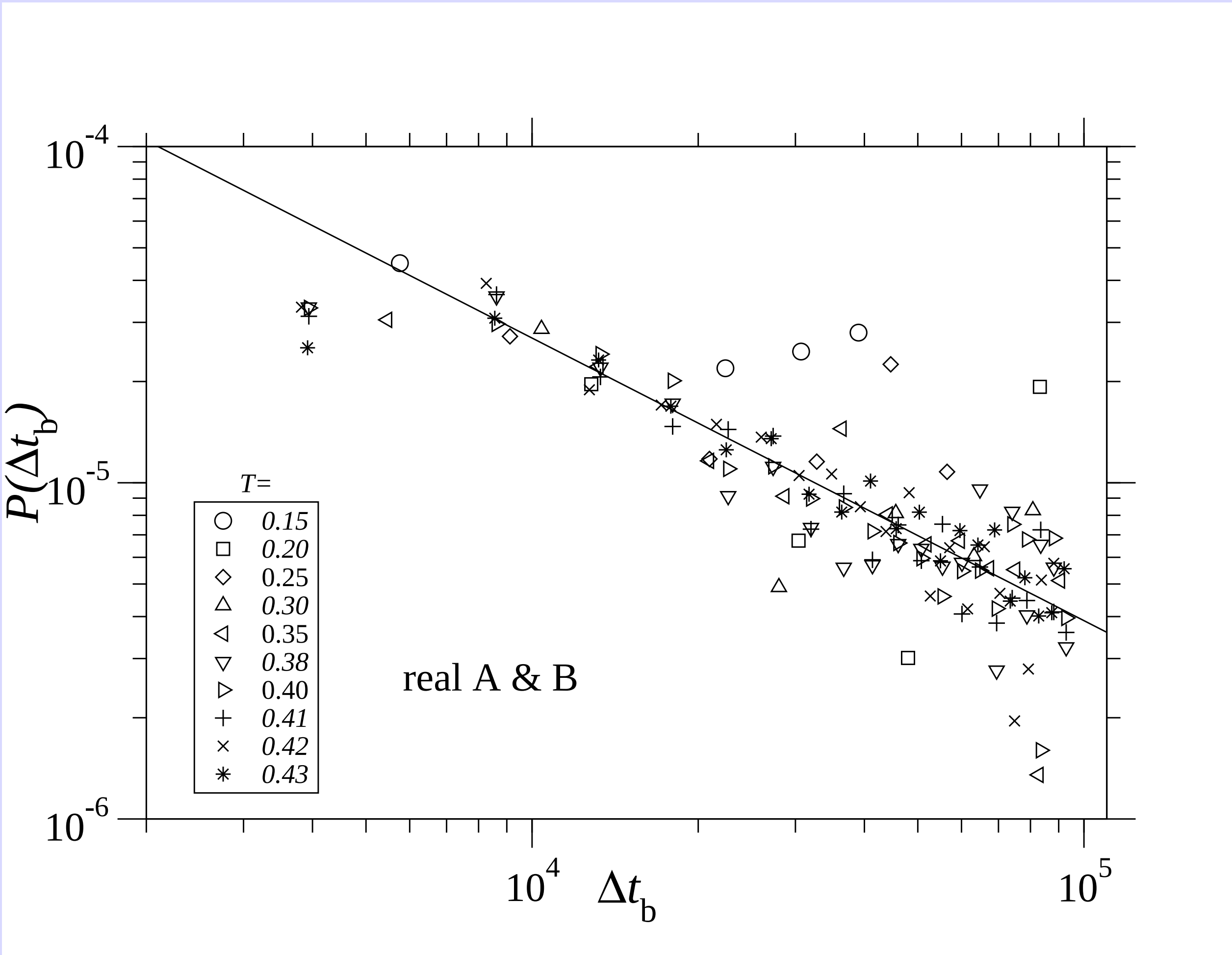
$$\vec{x}_i(t+\Delta t)=\vec{x}_i(t)+\vec{v}_i(t)\Delta t+\vec{a}_i(t)(\Delta t)^2/2$$

$$\vec{v}_i(t+\Delta t)=\vec{v}_i(t)+(\vec{a}_i(t)+\vec{a}_i(t+\Delta t))\,\Delta t/2$$

$$\vec{a}_i(t)=\vec{F}_i(t)/m_i = -\vec{\nabla}_i U(t)/m_i$$



slopes -6.3 for A and -4.8 for B particles $\longrightarrow$ subdiffusive



slopes -0.84 $\longrightarrow$ subdiffusive